

1 **MICROLOCAL PROPERTIES OF SEVEN-DIMENSIONAL LEMON AND**
2 **APPLE RADON TRANSFORMS WITH APPLICATIONS IN COMPTON**
3 **SCATTERING TOMOGRAPHY**

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ABSTRACT. We present a microlocal analysis of two novel Radon transforms of interest in Compton Scattering Tomography (CST), which map compactly supported L^2 functions to their integrals over seven-dimensional sets of apple and lemon surfaces. Specifically, we show that the apple and lemon transforms are elliptic Fourier Integral Operators (FIO), which satisfy the Bolker condition. After an analysis of the full seven-dimensional case, we focus our attention on n -D subsets of apple and lemon surfaces with fixed central axis, where $n < 7$. Such subsets of surface integrals have applications in airport baggage and security screening. When the data dimensionality is restricted, the apple transform is shown to violate the Bolker condition, and there are artifacts which occur on apple-cylinder intersections. The lemon transform is shown to satisfy the Bolker condition, when the support of the function is restricted to the strip $\{-1 < z < 1\}$.

6 1. INTRODUCTION

7 In this paper, we present a novel microlocal analysis of two Radon transforms of interest in
8 CST, which take the integrals of a function over seven dimensional sets of lemon and apple
9 surfaces. A “lemon” (also called a “spindle” in some works [25, 18, 24]) refers to the interior
10 part of a spindle (or self-intersecting) torus, and an “apple” is the exterior. See figure 1 for
11 a 2-D cross-section of a spindle torus, where we have highlighted the lemon and apple parts.
12 The literature considers lemon and apple transforms in 3-D CST [24, 25, 18, 23, 16, 17, 2],
13 where the goal is to reconstruct an electron density map from Compton scattered photons.
14 There is also a growing interest in the literature in Emission CST (ECST) [22, 9, 14, 13, 12],
15 where the aim is to reconstruct a gamma ray source from cone integral data.

16 In [17], two fixed-source CST configurations, with spherical and cylindrical detector ar-
17 rays, are considered. In both cases, the data is three dimensional, and consists of a two-
18 dimensional detector coordinate and a one-dimensional energy variable. Due to limited
19 energy resolution, the fixed source position, and the shape of the detector surface, the data
20 is incomplete. For example, the cylindrical acquisition geometry suffers limited angle issues.
21 In such cases of limited data, the reconstruction becomes unstable, and there are image
22 artifacts. The authors go on to develop a modified Kaczmarz algorithm to combat the re-
23 construction artifacts and test their algorithm on simulated examples with Poisson noise.

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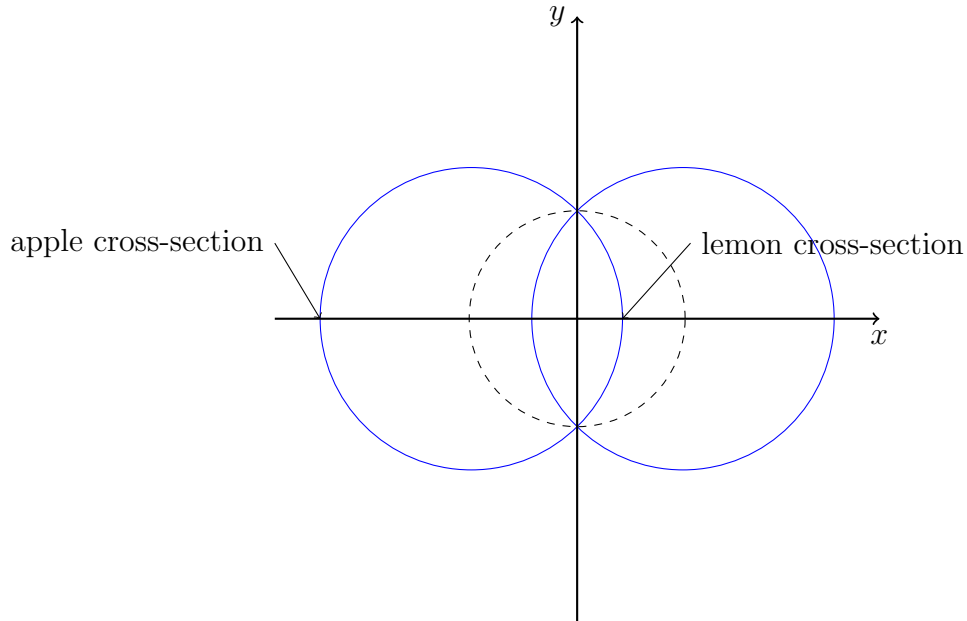


FIGURE 1. 2-D cross section of a spindle torus centered on the origin, with axis of revolution y . The lemon cross-section is the intersection of the interior of the dashed circle with the torus cross-section. The apple cross section is the intersection of the torus cross-section with the exterior of the dashed circle. The lemon/apple is the surface of revolution of the lemon/apple cross-section about y .

24 Similar reconstruction instabilities can be seen also in, e.g., conventional X-ray CT with
 25 limited angle data [1, 10, 11].

26 In [24], the authors present a microlocal analysis of the lemon transform introduced in [25].
 27 The acquisition geometry consists of a single rotating source and detector on a fixed axis. As
 28 in [17], the data is three-dimensional, and, in this case, consists of a 2-D rotation and a 1-D
 29 energy variable. The lemon transform is shown to violate the Bolker condition, and there are
 30 artifacts induced by flowout which appear as a spherical blurring effect in the reconstruction.
 31 There are also invisible singularities near the origin due to limited energy resolution. In [25],
 32 an algebraic reconstruction method is proposed to invert the lemon transform. Here artifacts
 33 are observed in reconstructions with noisy data, in line with the theory of [24].

34 In [2], the authors introduce a scanning modality in 3-D CST using a fixed source and
 35 single rotating detector restricted to a spherical surface. The data, in this case, has three
 36 degrees of freedom, and consists of a 2-D detector rotation and a 1-D energy variable. The
 37 authors model the Compton scatter intensity using a new apple Radon transform, and they
 38 derive an explicit inversion formula using a spherical harmonic expansion and Volterra inte-
 39 gral equation theory. Additionally, a hybrid analytic/algebraic reconstruction algorithm is
 40 presented and tested on simulated phantoms with added pseudo random noise. The authors
 41 discover blurring artifacts in the reconstructions, which indicate instabilities due to limited
 42 data, as is, for example, discovered in [17].

43 In the works discussed above, a number of imaging modalities are introduced based on
 44 practical machine designs, and the data dimension is such that the reconstruction target is

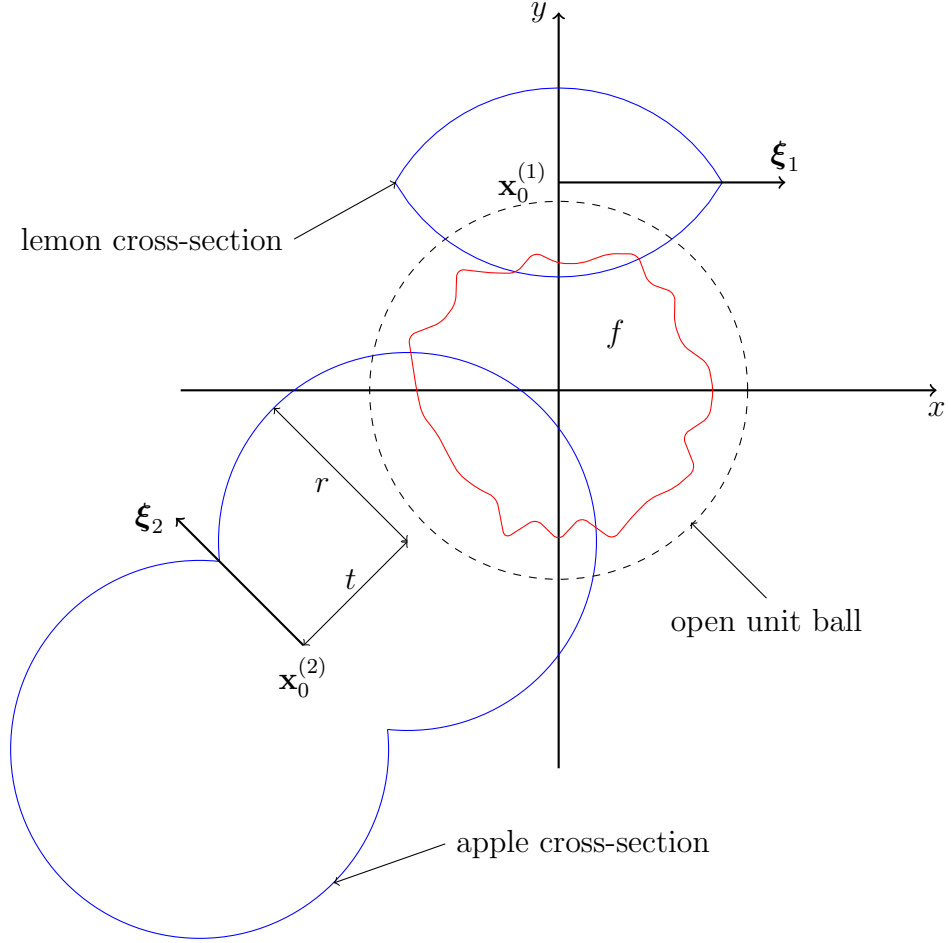


FIGURE 2

45 determined. That is, the reconstruction target and data are both three-dimensional. The set
 46 of spindle tori in 3-D space is seven-dimensional, and hence the literature thus far considers
 47 only limited data problems in CST, i.e., 3-D subsets of the full 7-D set of tori are considered.
 48 This often leads to artifacts and instabilities in the reconstruction due to, for example, limited
 49 angles (as in [17]) and failure to satisfy the Bolker condition [24]. In this paper, we wish to
 50 investigate the problem instability and presence of artifacts when there are no limits to the
 51 data dimensionality in CST, and we have knowledge of a seven-dimensional set of apple and
 52 lemon integrals in 3-D space. This can be considered a best case scenario in CST in terms of
 53 data dimensionality. Specifically, we consider the scanning geometry illustrated in figure 2.
 54 Here, we have shown an (x, y) plane cross-section of the scanning geometry. The scanning
 55 target (f) is supported on the open unit ball and is illustrated by an uneven red boundary.
 56 Example lemon and apple cross sections are drawn in blue, with centers $\mathbf{x}_0^{(1)}$ and $\mathbf{x}_0^{(2)}$, and
 57 axis of rotation ξ_1 and ξ_2 , respectively. The apple radius is denoted by r , and the distance
 58 from $\mathbf{x}_0^{(2)}$ to the center of the apple tube is denoted by t . We consider the apple and lemon
 59 surfaces whose points of self-intersection (which we will call *singular points*) lie outside the
 60 open unit ball. We do this to avoid singularities in the apple/lemon surface measure. In CST,
 61 the singular points of the lemons and apples correspond to source and detector coordinates.

62 So, in the context of CST, our geometry consists of all source and detector positions which
 63 lie outside the unit ball (this is a six-dimensional set). Additionally, we can vary the torus
 64 radius (r), which in CST is equivalent to the photon energy [17]. Thus, in total, our data
 65 set is seven-dimensional.

66 Motivated by the geometry of figure 2, we introduce novel lemon and apple Radon trans-
 67 forms, which map f to its integrals over seven-dimensional sets of apple and lemon surfaces.
 68 Our main theorem proves that the lemon and apple transforms are elliptic FIO which satisfy
 69 the Bolker condition. Additionally, we consider the practical applications of our theory to
 70 other scanning geometries from the literature. Specifically, we consider the scanning geome-
 71 try of [27], which is designed for use in airport baggage screening, and discuss the microlocal
 72 properties of lemon and apple transforms which induce translation on the scanning target.

73 The remainder of this paper is organized as follows. In section 2, we give some preliminary
 74 definitions and theorems that will be used in our analysis. In section 3, we introduce novel
 75 lemon and apple transforms, which map compactly supported L^2 functions to their integrals
 76 over seven-dimensional sets of lemon and apple surfaces, respectively, as pictured in figure
 77 2. Here we prove our main theorem, which shows that the lemon and apple transforms
 78 are elliptic FIO which satisfy the Bolker condition. In section 4, we consider a practical
 79 scanning geometry in CST, first introduced in [27], and discuss the artifacts in lemon and
 80 apple integral reconstructions when the axis of revolution of the lemons/apples is fixed, and
 81 the target function undergoes a 2-D translation.

82 2. DEFINITIONS AND PRELIMINARY THEOREMS

83 We next provide some notation and definitions. Let X and Y be open subsets of \mathbb{R}^n .
 84 Let $\mathcal{D}(X)$ be the space of smooth functions compactly supported on X with the standard
 85 topology and let $\mathcal{D}'(X)$ denote its dual space, the vector space of distributions on X . Let
 86 $\mathcal{E}(X)$ be the space of all smooth functions on X with the standard topology and let $\mathcal{E}'(X)$
 87 denote its dual space, the vector space of distributions with compact support contained in
 88 X . Finally, let $\mathcal{S}(\mathbb{R}^n)$ be the space of Schwartz functions, that are rapidly decreasing at ∞
 89 along with all derivatives. See [19] for more information.

90 For a function f in the Schwartz space $\mathcal{S}(\mathbb{R}^n)$ or in $L^2(\mathbb{R}^n)$, we use $\mathcal{F}f$ and $\mathcal{F}^{-1}f$ to denote
 91 the Fourier transform and inverse Fourier transform of f , respectively (see [6, Definition
 92 7.1.1]). Note that $\mathcal{F}^{-1}\mathcal{F}f(\mathbf{x}) = \frac{1}{(2\pi)^n} \int_{\mathbf{y} \in \mathbb{R}^n} \int_{\mathbf{z} \in \mathbb{R}^n} \exp((\mathbf{x} - \mathbf{z}) \cdot \mathbf{y}) f(\mathbf{z}) d\mathbf{z} d\mathbf{y}$.

93 We use the standard multi-index notation: if $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_n) \in \{0, 1, 2, \dots\}^n$ is a
 94 multi-index and f is a function on \mathbb{R}^n , then

$$\partial^\alpha f = \left(\frac{\partial}{\partial x_1} \right)^{\alpha_1} \left(\frac{\partial}{\partial x_2} \right)^{\alpha_2} \cdots \left(\frac{\partial}{\partial x_n} \right)^{\alpha_n} f.$$

95 If f is a function of $(\mathbf{y}, \mathbf{x}, \mathbf{s})$ then $\partial_{\mathbf{y}}^\alpha f$ and $\partial_{\mathbf{s}}^\alpha f$ are defined similarly.

96 We identify cotangent spaces on Euclidean spaces with the underlying Euclidean spaces,
 97 so we identify $T^*(X)$ with $X \times \mathbb{R}^n$.

98 If ϕ is a function of $(\mathbf{y}, \mathbf{x}, \mathbf{s}) \in Y \times X \times \mathbb{R}^N$ then we define $d_{\mathbf{y}}\phi = \left(\frac{\partial \phi}{\partial y_1}, \frac{\partial \phi}{\partial y_2}, \dots, \frac{\partial \phi}{\partial y_n} \right)$, and
 99 $d_{\mathbf{x}}\phi$ and $d_{\mathbf{s}}\phi$ are defined similarly. We let $d\phi = (d_{\mathbf{y}}\phi, d_{\mathbf{x}}\phi, d_{\mathbf{s}}\phi)$.

100 We use the convenient notation that if $A \subset \mathbb{R}^m$, then $\dot{A} = A \setminus \{\mathbf{0}\}$.

101 The singularities of a function and the directions in which they occur are described by the
 102 wavefront set [4, page 16]:

103 **Definition 2.1.** Let X Let an open subset of \mathbb{R}^n and let f be a distribution in $\mathcal{D}'(X)$. Let
 104 $(\mathbf{x}_0, \boldsymbol{\xi}_0) \in X \times \mathbb{R}^n$. Then f is *smooth at \mathbf{x}_0 in direction $\boldsymbol{\xi}_0$* if there exists a neighborhood U
 105 of \mathbf{x}_0 and V of $\boldsymbol{\xi}_0$ such that for every $\phi \in \mathcal{D}(U)$ and $N \in \mathbb{R}$ there exists a constant C_N such
 106 that for all $\boldsymbol{\xi} \in V$,

$$(2.1) \quad |\mathcal{F}(\phi f)(\lambda \boldsymbol{\xi})| \leq C_N(1 + |\lambda|)^{-N}.$$

107 The pair $(\mathbf{x}_0, \boldsymbol{\xi}_0)$ is in the *wavefront set*, $\text{WF}(f)$, if f is not smooth at \mathbf{x}_0 in direction $\boldsymbol{\xi}_0$.

108 This definition follows the intuitive idea that the elements of $\text{WF}(f)$ are the point–normal
 109 vector pairs above points of X at which f has singularities. For example, if f is the char-
 110 acteristic function of the unit ball in \mathbb{R}^3 , then its wavefront set is $\text{WF}(f) = \{(\mathbf{x}, t\mathbf{x}) : \mathbf{x} \in$
 111 $S^2, t \neq 0\}$, the set of points on a sphere paired with the corresponding normal vectors to the
 112 sphere.

113 The wavefront set of a distribution on X is normally defined as a subset the cotangent
 114 bundle $T^*(X)$ so it is invariant under diffeomorphisms, but we do not need this invariance,
 115 so we will continue to identify $T^*(X) = X \times \mathbb{R}^n$ and consider $\text{WF}(f)$ as a subset of $X \times \mathbb{R}^n$.

116 **Definition 2.2** ([6, Definition 7.8.1]). We define $S^m(Y \times X \times \mathbb{R}^N)$ to be the set of $a \in$
 117 $\mathcal{E}(Y \times X \times \mathbb{R}^N)$ such that for every compact set $K \subset Y \times X$ and all multi–indices α, β, γ
 118 the bound

$$|\partial_{\mathbf{y}}^{\gamma} \partial_{\mathbf{x}}^{\beta} \partial_{\boldsymbol{\sigma}}^{\alpha} a(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma})| \leq C_{K, \alpha, \beta, \gamma} (1 + \|\boldsymbol{\sigma}\|)^{m - |\alpha|}, \quad (\mathbf{y}, \mathbf{x}) \in K, \boldsymbol{\sigma} \in \mathbb{R}^N,$$

119 holds for some constant $C_{K, \alpha, \beta, \gamma} > 0$.

120 The elements of S^m are called *symbols* of order m . Note that these symbols are sometimes
 121 denoted $S_{1,0}^m$. The symbol $a \in S^m(Y, X, \mathbb{R}^N)$ is *elliptic* if for each compact set $K \subset Y \times X$,
 122 there is a $C_K > 0$ and $M > 0$ such that

$$(2.2) \quad |a(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma})| \geq C_K (1 + \|\boldsymbol{\sigma}\|)^m, \quad (\mathbf{y}, \mathbf{x}) \in K, \|\boldsymbol{\sigma}\| \geq M.$$

123 **Definition 2.3** ([7, Definition 21.2.15]). A function $\phi = \phi(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma}) \in \mathcal{E}(Y \times X \times \mathbb{R}^N)$ is a
 124 *phase function* if $\phi(\mathbf{y}, \mathbf{x}, \lambda \boldsymbol{\sigma}) = \lambda \phi(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma})$, $\forall \lambda > 0$ and $d\phi$ is nowhere zero. The *critical set*
 125 *of ϕ* is

$$\Sigma_{\phi} = \{(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma}) \in Y \times X \times \mathbb{R}^N : d_{\boldsymbol{\sigma}} \phi = 0\}.$$

126 A phase function is *clean* if the critical set $\Sigma_{\phi} = \{(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma}) : d_{\boldsymbol{\sigma}} \phi(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma}) = 0\}$ is a smooth
 127 manifold with tangent space defined as the kernel of $d(d_{\boldsymbol{\sigma}} \phi)$ on Σ_{ϕ} . Here, the derivative
 128 d is applied component–wise to the vector–valued function $d_{\boldsymbol{\sigma}} \phi$. So, $d(d_{\boldsymbol{\sigma}} \phi)$ is treated as a
 129 Jacobian matrix of dimensions $N \times (2n + N)$.

130 By the Constant Rank Theorem the requirement for a phase function to be clean is satisfied
 131 if $d(d_{\boldsymbol{\sigma}} \phi)$ has constant rank.

132 **Definition 2.4** ([7, Definition 21.2.15] and [8, section 25.2]). Let X and Y be open subsets
 133 of \mathbb{R}^n . Let $\phi \in \mathcal{E}(Y \times X \times \mathbb{R}^N)$ be a clean phase function. In addition, we assume that ϕ
 134 is *nondegenerate* in the following sense:

$$d_{\mathbf{y}} \phi \text{ and } d_{\mathbf{x}} \phi \text{ are never zero on } \Sigma_{\phi}.$$

135 The *canonical relation parametrized by ϕ* is defined as

$$(2.3) \quad \mathcal{C} = \{((\mathbf{y}, d_{\mathbf{y}} \phi(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma})); (\mathbf{x}, -d_{\mathbf{x}} \phi(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma}))) : (\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma}) \in \Sigma_{\phi}\},$$

136 **Definition 2.5.** Let X and Y be open subsets of \mathbb{R}^n . Let an operator $A : \mathcal{D}(X) \rightarrow$
137 $\mathcal{D}'(Y)$ be defined by the distribution kernel $K_A \in \mathcal{D}'(X \times Y)$, in the sense that $Af(\mathbf{y}) =$
138 $\int_X K_A(\mathbf{x}, \mathbf{y})f(\mathbf{x})d\mathbf{x}$. Then we call K_A the *Schwartz kernel* of A . A *Fourier integral operator*
139 *(FIO)* of order $m + N/2 - n/2$ is an operator $A : \mathcal{D}(X) \rightarrow \mathcal{D}'(Y)$ with Schwartz kernel given
140 by an oscillatory integral of the form

$$(2.4) \quad K_A(\mathbf{y}, \mathbf{x}) = \int_{\mathbb{R}^N} e^{i\phi(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma})} a(\mathbf{y}, \mathbf{x}, \boldsymbol{\sigma}) d\boldsymbol{\sigma},$$

141 where ϕ is a clean nondegenerate phase function and a is a symbol in $S^m(Y \times X \times \mathbb{R}^N)$. The
142 *canonical relation* of A is the canonical relation of ϕ defined in (2.3).

143 The FIO A is *elliptic* if its symbol is elliptic.

144 This is a simplified version of the definition of FIO in [3, section 2.4] or [8, section 25.2]
145 that is suitable for our purposes since our phase functions are global. Because we assume
146 phase functions are nondegenerate, our FIO can be defined as maps from $\mathcal{E}'(X)$ to $\mathcal{D}'(Y)$
147 and sometimes on larger domains. For general information about FIOs, see [3, 8, 7]. For
148 information about the Schwartz Kernel, see [6, Theorem 5.1.9].

149 .

150 Let X and Y be sets and let $\Omega_1 \subset X$ and $\Omega_2 \subset Y \times X$. The composition $\Omega_2 \circ \Omega_1$ and
151 transpose Ω_2^t of Ω_2 are defined

$$\begin{aligned} \Omega_2 \circ \Omega_1 &= \{y \in Y : \exists x \in \Omega_1, (y, x) \in \Omega_2\} \\ \Omega_2^t &= \{(x, y) : (y, x) \in \Omega_2\}. \end{aligned}$$

152

153 The Hörmander-Sato Lemma provides the relationship between the wavefront set of dis-
154 tributions and their images under FIO.

155 **Theorem 2.6** ([6, Theorem 8.2.13]). *Let $f \in \mathcal{E}'(X)$ and let $A : \mathcal{E}'(X) \rightarrow \mathcal{D}'(Y)$ be an FIO*
156 *with canonical relation \mathcal{C} . Then, $\text{WF}(Af) \subset \mathcal{C} \circ \text{WF}(f)$.*

157 **Definition 2.7.** Let $\mathcal{C} \subset T^*(Y \times X)$ be the canonical relation associated to the FIO $A :$
158 $\mathcal{E}'(X) \rightarrow \mathcal{D}'(Y)$. We let Π_L and Π_R denote the natural left- and right-projections of \mathcal{C} ,
159 projecting onto the appropriate coordinates: $\Pi_L : \mathcal{C} \rightarrow T^*(Y)$ and $\Pi_R : \mathcal{C} \rightarrow T^*(X)$.

160 Because ϕ is nondegenerate, the projections do not map to the zero section.

161 Let A be an FIO with adjoint A^* . If A satisfies our next definition, then A^*A (or, if A does
162 not map to $\mathcal{E}'(Y)$, then $A^*\psi A$ for an appropriate cutoff ψ) is a pseudodifferential operator
163 [5, 15].

164 **Definition 2.8.** Let $A : \mathcal{E}'(X) \rightarrow \mathcal{D}'(Y)$ be a FIO with canonical relation \mathcal{C} then A (or \mathcal{C})
165 satisfies the *semi-global Bolker Condition* if the natural projection $\Pi_L : \mathcal{C} \rightarrow T^*(Y)$ is an
166 embedding (injective immersion).

Theorem 2.9 (Sylvester's Determinant Theorem (SDT) [20, 21]). *Let A be an $m \times n$ matrix,*
and B an $n \times m$ matrix. Then

$$\det(I_{m \times m} + AB) = \det(I_{n \times n} + BA).$$

168 In this section, we present a microlocal analysis of two new Radon transforms which map
 169 compactly supported L^2 functions to their integrals over seven-dimensional sets of lemon
 170 and apple surfaces. First, we give the defining equations for the apple and lemon surfaces.

171 Spindle tori are described by their center, $\mathbf{x}_0 \in \mathbb{R}^3$, their axis of revolution, and parameters
 172 s and t ; \sqrt{s} is the radius and t is the tube radius of the spindle torus. If ℓ is a line through
 173 the origin parallel to the axis of revolution of a spindle torus, then for some $\omega \in S^2$, one can
 174 write

$$\ell = \mathbb{R}\omega := \{\nu\omega : \nu \in \mathbb{R}\}$$

175 and the axis of revolution of the torus is $\mathbf{x}_0 + \ell$. We will call this line ℓ the *directional axis*
 176 of the spindle torus (equivalently, of the apple or lemon).

177 We will use rotation matrices to describe the directional axes of spindle tori. Let $(\alpha, \beta) \in$
 178 $[0, 2\pi] \times [0, \pi/2]$. Then, we define

$$(3.1) \quad \begin{aligned} R = R(\alpha, \beta) &= \begin{pmatrix} \cos \alpha & -\sin \alpha & 0 \\ \sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \beta & -\sin \beta \\ 0 & \sin \beta & \cos \beta \end{pmatrix} \\ &= \begin{pmatrix} \cos \alpha & -\sin \alpha \cos \beta & \sin \alpha \sin \beta \\ \sin \alpha & \cos \alpha \cos \beta & -\cos \alpha \sin \beta \\ 0 & \sin \beta & \cos \beta \end{pmatrix}. \end{aligned}$$

179 Let

$$(3.2) \quad \mathbf{x} = (x, y, z), \quad \mathbf{x}_0 = (x_0, y_0, z_0), \quad \mathbf{x}_T = \mathbf{x} - \mathbf{x}_0, \quad \mathbf{x}' = (x', y', z') = R^T(\alpha, \beta)\mathbf{x}_T$$

180 and

$$(3.3) \quad h(t, \mathbf{x}_0; \mathbf{x}) = \|\mathbf{x}_T\|^2 + t^2, \quad g(\alpha, \mathbf{x}_0, \beta; \mathbf{x}) = \sqrt{x'^2 + y'^2}.$$

181 We now define

$$(3.4) \quad \begin{aligned} \Psi_j(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}) &:= \left(\sqrt{x'^2 + y'^2} + (-1)^j t \right)^2 + z'^2 - s \\ &= h(t, \mathbf{x}_0; \mathbf{x}) + 2t(-1)^j g(\mathbf{x}_0, \alpha, \beta; \mathbf{x}) - s. \end{aligned}$$

182 and

$$(3.5) \quad T_j(s, t, \mathbf{x}_0, \alpha, \beta) = \{\mathbf{x} \in B : \Psi_j(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}) = 0\},$$

183 for $j = 1, 2$. Ψ_1 and Ψ_2 are the defining equations for apple and lemon surfaces, respectively,
 184 and T_1 and T_2 are the intersections of apples and lemons with B . See figure 3 for example
 185 2-D cross sections of apples and lemons with the defining equations highlighted.

186 **3.1. Definition of apple and lemon transforms.** Throughout this paper, we let $L_c^2(X)$
 187 denote the set of L^2 functions compactly supported on $X \subset \mathbb{R}^3$. Recall that the two points
 188 of intersection of the apple (resp. lemon) with its axis of revolution are called the *singular*
 189 *points* of the apple (resp. lemon). Note that the singular points are the points of intersection
 190 of the apple and lemon with the same parameters in Y , and they are singular points of both
 191 the apple and lemon. We will define the apple and lemon transforms on functions $f \in L_c^2(X)$,

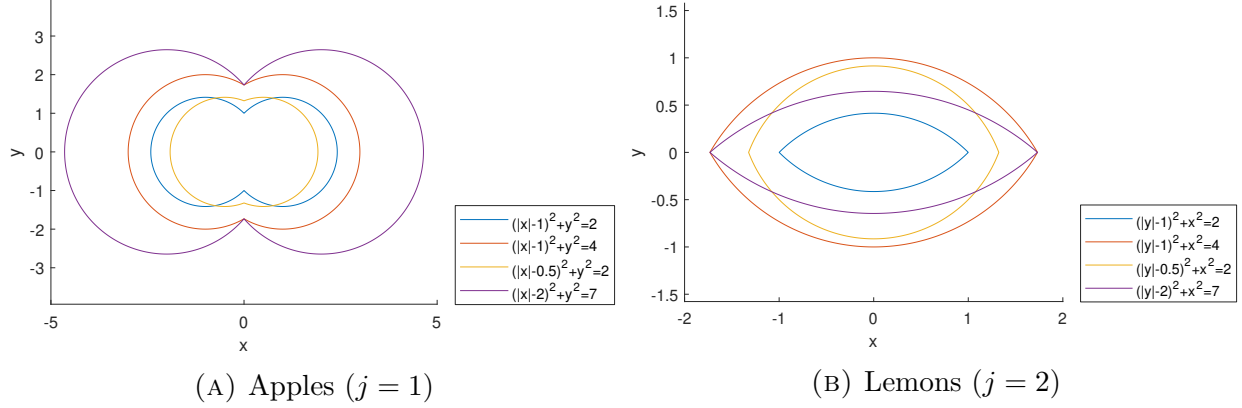


FIGURE 3. (x, y) plane cross sections of the apple and lemon parts of a spindle torus when $R = R(0, \frac{\pi}{2})$ (left) and $R = R(\frac{\pi}{2}, \frac{\pi}{2})$ (right), $\mathbf{x}_0 = \mathbf{0}$, and s and t vary between $\frac{1}{2}$ and 7.

192 where $X = B$ is the open unit ball in \mathbb{R}^3 , and we will need to ensure that the singular points
 193 of the apple or lemon do not meet the closed unit ball, \overline{B} . For this reason, we define

$$(3.6) \quad Y = \{(s, t, \mathbf{x}_0, \alpha, \beta) \in \mathbb{R}^2 \times \mathbb{R}^3 \times [0, 2\pi] \times [0, \pi/2] \\ : s > t^2, \{\mathbf{x}_0 \pm \sqrt{s - t^2} R(\alpha, \beta) \mathbf{e}_3\} \cap \overline{B} = \emptyset\},$$

194 where $\mathbf{e}_3 = (0, 0, 1)^T$ is the north pole. Note that every apple (for $j = 1$) and lemon
 195 (for $j = 2$) with singular points not meeting \overline{B} can be written $T_j(s, t, \mathbf{x}_0, \alpha, \beta)$ for some
 196 $(s, t, \mathbf{x}_0, \alpha, \beta) \in Y$ because all directional axes are generated by the map

$$(3.7) \quad [0, 2\pi] \times [0, \pi/2] \ni (\alpha, \beta) \mapsto \mathbb{R}R(\alpha, \beta)\mathbf{e}_3.$$

197 *Remark 3.1.* This map (3.7) from (α, β) to directional axes is not injective for $\beta = 0, \pi/2$.
 198 Therefore, we cannot use $[0, 2\pi] \times [0, \pi/2]$ to parameterize direction axes, as it would cause
 199 issue later in the proofs of our main theorems. Furthermore, our parameter space Y in (3.6)
 200 is not a manifold without boundary because $[0, \pi/2]$ is not a manifold without boundary.
 201 Note that we are identifying 0 and 2π to transform $[0, 2\pi]$ to the manifold S^1 .

202 At the start of the proof of Theorem 3.2, we will define a parameter set for spindle tori
 203 that is a manifold without boundary for which the map to spindle tori (with singular points
 204 outside \overline{B}) is bijective. These properties are required to use the standard microlocal analysis
 205 of Radon transforms (e.g., see [5]). However, we will parameterize spindle tori using Y when
 206 appropriate.

207 We define the Radon transforms which take the integrals of f over apple ($j = 1$) and
 208 lemon ($j = 2$) surfaces

$$(3.8) \quad \mathcal{R}_j f(s, t, \mathbf{x}_0, \alpha, \beta) = \int_X \|\nabla_{\mathbf{x}} \Psi_j\| \delta(\Psi_j(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x})) f(\mathbf{x}) d\mathbf{x} \\ = \int_{-\infty}^{\infty} \int_X \|\nabla_{\mathbf{x}} \Psi_j\| e^{\sigma \Psi_j(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x})} f(\mathbf{x}) d\mathbf{x} d\sigma,$$

and we let

$$\mathcal{A}f = \mathcal{R}_1 f, \quad \mathcal{L}f = \mathcal{R}_2 f$$

209 where \mathcal{A} is called the *apple transform*, and \mathcal{L} is the *lemon transform*.

210 Here, we will assume the gradient of a scalar valued function is a column vector, as are
211 elements of \mathbb{R}^n .

212 To ensure that $T_j(s, t, \mathbf{x}_0, \alpha, \beta)$ is a smooth manifold and that the weight $\|\nabla_{\mathbf{x}}\Psi_j\|$ in (3.8)
213 is defined, we have defined Y so that it includes only the apple and lemon surfaces whose
214 singular points do not intersect $X = B$. This way, in the integrals of (3.8), we stay away from
215 the singular points of the apples and lemons, and any singularities in the FIO amplitudes
216 and phases. Strictly speaking, one would add a smooth cutoff to the symbol which is zero
217 close to the central axis of the spindle tori, as in [26, Lemma 3.3], so the amplitude is smooth
218 everywhere and the phase is smooth near the support of the amplitude. However, we do not
219 go into such technicalities here.

220 Now that the apple and lemon transforms are defined we present a separate microlocal
221 analysis of each transform in the following sections.

222 **3.2. Microlocal properties of \mathcal{A} ; the $j = 1$ case.** Here we discuss the microlocal prop-
223 erties of the apple transform \mathcal{A} . Our first theorem proves that \mathcal{A} is an elliptic FIO.

224 **Theorem 3.2.** *The apple transform $\mathcal{A} = \mathcal{R}_1$ of (3.8) is an elliptic FIO order -2 from*
225 *domain $\mathcal{E}'(B)$ to $\mathcal{D}'(Y)$.*

226 *Proof.* To analyze \mathcal{A} as an FIO, we need to parametrize apples using a manifold without
227 boundary, as discussed in Remark 3.1. However, Y cannot be used, since it is not a manifold
228 without boundary since $[0, \pi/2]$ has boundary points $0, \pi/2$. To get around this, we first
229 parametrize all spindle tori in a global way as a manifold without boundary. This is required
230 to use the theory of Radon transforms as FIO [5]. To define this manifold, we parametrize
231 spindle tori by points $\mathbf{y} = (s, t, \mathbf{x}_0, \ell)$, as discussed at the start of this section, where \sqrt{s} is
232 the radius, t is the tube radius of the spindle torus, \mathbf{x}_0 is its center, and ℓ is the directional
233 axis. Recall that the directional axis of a spindle torus is the line through the origin in \mathbb{R}^3 ,
234 which is parallel to the axis of revolution of the torus, $\mathbf{x}_0 + \ell$. The set of lines through the
235 origin in \mathbb{R}^3 is denoted \mathbb{RP}^2 and is called the two-dimensional real projective space.

236 We let \tilde{Y} be the set of $\mathbf{y} = (s, t, \mathbf{x}_0, \ell)$ such that the singular points of the spindle torus
237 parameterized by \mathbf{y} do not meet \bar{B} . Then, \tilde{Y} is a manifold without boundary that parame-
238 terizes all apples ($j = 1$) and all lemons ($j = 2$) the singular points of which do not meet \bar{B}
239 by the map

$$(3.9) \quad \begin{aligned} \tilde{Y} \ni (s, t, \mathbf{x}_0, \ell) &\mapsto T_j(s, t, \mathbf{x}_0, \alpha, \beta) \\ &\text{when } (\alpha, \beta) \text{ is chosen so that} \\ &\ell = \mathbb{R}R(\alpha, \beta)\mathbf{e}_3. \end{aligned}$$

240 Note that the map in (3.9) and $\mathcal{A}f$ are well-defined on \tilde{Y} because the spindle torus and
241 its measure are the same no matter which (α, β) one chooses that satisfies $\ell = \mathbb{R}R(\alpha, \beta)\mathbf{e}_3$.
242 This is true by rotation invariance of the spindle torus about its axis of revolution, $\mathbf{x}_0 + \ell$
243 and rotation invariance of the integral over the torus. Furthermore, every spindle torus is
244 described by a unique $(s, t, \mathbf{x}_0, \ell)$.

245 To get local coordinates on \tilde{Y} , we need to specify local coordinates on \mathbb{RP}^2 , since (s, t, \mathbf{x}_0)
246 are already coordinates. We choose a vertical axis and let \mathbf{e}_3 be the unit vector pointing in
247 the positive direction along that axis. Then, we let \mathbf{e}_1 and \mathbf{e}_2 be orthogonal unit vectors so

248 $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ form a right-hand coordinate system in \mathbb{R}^3 . Now, we define the domain of the
 249 coordinate map

$$(3.10) \quad Y' = \{(s, t, \mathbf{x}_0, \alpha, \beta) \in Y : \beta \in (0, \pi/2)\}.$$

250 Then, local coordinates on \tilde{Y} are given by

$$(3.11) \quad Y' \ni (s, t, \mathbf{x}_0, \alpha, \beta) \mapsto (s, t, \mathbf{x}_0, \mathbb{R}R(\alpha, \beta)\mathbf{e}_3).$$

251 For different choices of basis $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ on \mathbb{R}^3 with vertical axis in direction of \mathbf{e}_3 , this
 252 coordinate map describes a coordinate chart on \tilde{Y} .

253 We will work in these coordinates and use the notation (3.1)-(3.4), and (3.8) for the rest
 254 of this section.

From (3.8), the phase function of \mathcal{R}_1 is

$$\Phi_1(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}; \sigma) = \sigma \Psi_1(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}).$$

255 We now show that Φ_1 is clean, non-degenerate and homogeneous in σ order 1, so that \mathcal{A}
 256 satisfies the definition of FIO (see definition 2.5). Φ_1 is trivially homogeneous order 1, since
 257 Ψ_1 does not depend on σ . Note also, $d_s \Phi_1 = -\sigma \neq 0$, hence $d\Phi_1, d_y \Phi_1 \neq 0$. The apple
 258 surfaces are smooth manifolds away from their singular points—the points which we do not
 259 consider. Hence Φ_1 is clean.

Let $\mathbf{x}_0 = (x_0, y_0, z_0)$, then we will let x'_{x_0} denote the partial derivative of $x' = x'(\mathbf{x}_0, \alpha, \beta; \mathbf{x})$
 with respect to x_0 and define the other partial derivatives of x', y', z' analogously. Let

$$R^T = \begin{pmatrix} \cos \alpha & \sin \alpha & 0 \\ -\sin \alpha \cos \beta & \cos \alpha \cos \beta & \sin \beta \\ \sin \alpha \sin \beta & -\cos \alpha \sin \beta & \cos \beta \end{pmatrix} = \begin{pmatrix} r_1 \\ r_2 \\ r_3 \end{pmatrix}$$

260 have rows r_1, r_2, r_3 . Then, we have

$$(3.12) \quad \begin{aligned} \nabla_{\mathbf{x}_0} \Phi_1 &= -2\sigma \left[\mathbf{x}_T + \frac{t}{g} (x'_{x_0} x' + y'_{x_0} y', x'_{y_0} x' + y'_{y_0} y', x'_{z_0} x' + y'_{z_0} y')^T \right] \\ &= -2\sigma \left(\mathbf{x}_T + \frac{t}{g} (\nabla_{\mathbf{x}_0} x', \nabla_{\mathbf{x}_0} y') \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T \right) \\ &= -2\sigma \left(I - \frac{t}{g} A \right) \mathbf{x}_T \\ &= -\sigma \nabla_{\mathbf{x}} \Psi_1 \\ &= -d_{\mathbf{x}} \Phi_1, \end{aligned}$$

261 where

$$(3.13) \quad A = (r_1^T, r_2^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix}$$

262 is symmetric, idempotent (i.e., $A^T = A$ and $A^2 = A$) and , and I is the 3×3 identity matrix.
 263 We have

$$\begin{aligned}
 \det \left(I - \frac{t}{g} A \right) &= \det \left(I - \frac{t}{g} (r_1^T, r_2^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \right) \\
 &= \det \left(I_{2 \times 2} - \frac{t}{g} \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} (r_1^T, r_2^T) \right), \quad (\text{by SDT}) \\
 (3.14) \qquad &= \det \left(I_{2 \times 2} - \frac{t}{g} I_{2 \times 2} \right) \\
 &= \left(1 - \frac{t}{g} \right)^2,
 \end{aligned}$$

264 which is zero if and only if $t = g$. Recall that we exclude the case $g = 0$ since the singular
 265 points of apples parameterized by Y or \tilde{Y} do not meet \bar{B} .

266 If $t \neq g$, then $\left(I - \frac{t}{g} A \right)$ is invertible and $\left(I - \frac{t}{g} A \right) \mathbf{x}_T = 0 \implies \mathbf{x}_T = 0$ but this would
 267 mean that the center of the apple, \mathbf{x}_0 , is on the apple (equivalently, $t^2 - s = 0$). However,
 268 $s > t^2 > 0$ so this is not possible.

269 Now, we consider the case when $t = g$. Let C_t denote the cylinder of radius t with axis of
 270 revolution $\{\mathbf{x}_0 + pR\mathbf{e}_3 : p \in \mathbb{R}\}$. If $t = g$, then \mathbf{x} is in C_t and in $\text{Null}(I - A) = \text{span}(r_1^T, r_2^T)$.
 271 Thus, if $t = g$, and $(I - A)\mathbf{x}_T = 0$, then $\mathbf{x} \in \{\mathbf{x}_0 + \text{span}(r_1^T, r_2^T)\} \cap C_t$. If \mathbf{x} is in the critical
 272 set of Φ_1 also (i.e., \mathbf{x} lies on the torus parameterized by \mathbf{y}) then s must be zero (i.e., the
 273 apple radius is zero), which we do not consider since $\sqrt{s} > t > 0$. Therefore, $d_{\mathbf{x}}\Phi_1 \neq 0$ and
 274 Φ_1 is nondegenerate.

275 The amplitude of \mathcal{A} is

$$\begin{aligned}
 a_1(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}) &= \|\nabla_{\mathbf{x}}\Psi_1(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x})\| \\
 (3.15) \qquad &= 2 \left(\mathbf{x}_T^T \left(I - \frac{t}{g} A \right)^T \left(I - \frac{t}{g} A \right) \mathbf{x}_T \right)^{\frac{1}{2}}
 \end{aligned}$$

276 by (3.12). By the arguments of the last paragraph we can show that a_1 is never zero, and
 277 hence a_1 is an elliptic symbol. a is order zero since it is smooth, and does not depend on σ .
 278 Hence, \mathcal{A} is an elliptic FIO order $O(\mathcal{A}) = 0 + \frac{1}{2} - \frac{7}{2} = -2$. \square

279 We now have our first main theorem which shows that \mathcal{A} satisfies the semiglobal Bolker
 280 condition.

281 **Theorem 3.3.** *The left projection $\Pi_L^{(1)}$ of \mathcal{A} is an injective immersion, and hence \mathcal{A} satisfies*
 282 *the semiglobal Bolker condition from domain $\mathcal{E}'(B)$ to $\mathcal{D}'(Y)$.*

283 As the proof for Theorem 3.3 is long, we split the proof into two subsections. We start
 284 with the immersion proof in the next section, and present proof of injectivity in the following
 285 section.

286 3.2.1. $\Pi_L^{(1)}$ immersion proof. Since being an immersion is a local property , we can check this
 287 at an arbitrary point $(\mathbf{y}, \eta, \mathbf{x}, \xi)$ in the canonical relation of \mathcal{A} . Let ℓ be the direction axis of
 288 revolution of the apple parameterized by \mathbf{y} . Choose a unit vector \mathbf{e}_3 such that ℓ is neither
 289 parallel nor perpendicular to \mathbf{e}_3 . Choose unit vectors \mathbf{e}_1 and \mathbf{e}_2 so that $(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ makes

290 up a right-hand coordinate system on \mathbb{R}^3 . Use this coordinate system on \mathbb{R}^3 to define the
 291 coordinate map (3.11) and the set Y' (see (3.10)). Throughout this proof, the calculations
 292 are performed using this coordinate system.

First we calculate $\Pi_L^{(1)}$. We have the derivatives

$$293 \quad d_t \Phi_1 = 2\sigma(t - g), \quad d_s \Phi_1 = -\sigma,$$

$$(3.16) \quad d_\alpha \Phi_1 = \frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T,$$

294 and

$$(3.17) \quad d_\beta \Phi_1 = \frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\beta}^T, r_{2\beta}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T,$$

295 where $r_{i\alpha}$ is the component-wise partial derivative of r_i with respect to α (similarly for $r_{i\beta}$).
 296 Let

$$H = \{\mathbf{x} \in \mathbb{R}^3 : \Psi_1(t^2, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}) = 0\}$$

297 be the horn torus with radius t and axis of revolution $\mathbf{x}_0 + \mathbb{R}R(\alpha, \beta)\mathbf{e}_3$, and let $\tilde{H} = \{\mathbf{x} \in \mathbb{R}^3 : \Psi_1(t^2, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}) > 0\}$ be the exterior of H . Let $\mathcal{D}_1 = \{(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}; \sigma) : \mathbf{x} \in \tilde{H}\} \times \mathbb{R} \setminus \{0\}$,
 298 then the map
 299

$$\mathcal{D}_1 \ni (t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}; \sigma) \mapsto (s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}; d_s \Phi_1, d_t \Phi_1, d_\beta \Phi_1; \nabla_{\mathbf{x}_0} \Phi_1; \nabla_{\mathbf{x}} \Phi_1) \in \mathcal{C}_1$$

300 where $s = h - 2tg$ gives local coordinates on the canonical relation for $\mathcal{A} = \mathcal{R}_1$.

301 In these coordinates, the left projection $\Pi_L^{(1)} : \mathcal{D}_1 \rightarrow \Pi_L^{(1)}(\mathcal{D}_1)$ of \mathcal{A} is defined

$$(3.18) \quad \Pi_L^{(1)}(\sigma; t, \alpha, \beta, \mathbf{x}_0; \mathbf{x}) = \left(\underbrace{-\sigma}_{d_s \Phi_1}, t, \alpha, \beta, \mathbf{x}_0, \overbrace{-2\sigma \mathbf{x}_T^T \left(I - \frac{t}{g} A \right)}^{\nabla_{\mathbf{x}_0} \Phi_1}, \underbrace{2\sigma(t - g)}_{d_t \Phi_1}, \right. \\ \left. \underbrace{\frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T}_{d_\alpha \Phi_1}, \underbrace{\frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\beta}^T, r_{2\beta}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T}_{d_\beta \Phi_1}, \underbrace{h - 2tg}_s \right),$$

302 where we have highlighted the derivatives of Φ_1 using under and overbraces. Also, we have
 303 rearranged the variables in (3.18) to correspond to the order used in calculating the Jacobian
 304 matrix of $\Pi_L^{(1)}$:

$$(3.19) \quad D\Pi_L^{(1)} = \begin{matrix} & & d\sigma, dt, d\alpha, d\beta, D\mathbf{x}_0 & & \nabla_{\mathbf{x}} \\ & & & & \mathbf{0}_{7 \times 3} \\ d_s \Phi_1, t, \alpha, \beta, \mathbf{x}_0 & \left(\begin{array}{c} I'_{7 \times 7} \\ \cdot \\ \cdot \\ \cdot \\ \cdot \\ \cdot \\ \cdot \end{array} \right. & & & \\ \nabla_{\mathbf{x}_0} \Phi_1 & & & & D_{\mathbf{x}} \left(-2\sigma \mathbf{x}_T^T \left(I - \frac{t}{g} A \right) \right) \\ d_t \Phi_1 & & & & \nabla_{\mathbf{x}} (2\sigma(t - g))^T \\ d_\alpha \Phi_1 & & & & \nabla_{\mathbf{x}} \left(\frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T \right)^T \\ d_\beta \Phi_1 & & & & \nabla_{\mathbf{x}} \left(\frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\beta}^T, r_{2\beta}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T \right)^T \\ s & & & & \nabla_{\mathbf{x}} (h - 2tg)^T \end{matrix} \right),$$

305 where $I'_{n \times n}$ is the $n \times n$ identity matrix but with the first entry replaced by -1 . Here we
306 have highlighted the arguments of $\Pi_L^{(1)}$ on the left-hand side of the matrix for $D\Pi_L^{(1)}$, and
307 the order of derivatives is indicated above $D\Pi_L^{(1)}$. The terms corresponding to \cdot in $D\Pi_L^{(1)}$ are
308 not important for our calculations, as they will be multiplied by zero in the calculation of
309 the determinant of $D\Pi_L^{(1)}$. We now find the derivatives in the right-hand column of $D\Pi_L^{(1)}$
310 and show that $D\Pi_L^{(1)}$ is full rank.

Using the product rule and

$$\nabla_{\mathbf{x}} \left(\frac{1}{g} \right) = -\frac{1}{g^3} A \mathbf{x}_T,$$

311 we can calculate the Jacobian matrix

$$(3.20) \quad \begin{aligned} D_{\mathbf{x}} \left(-2\sigma \mathbf{x}_T^T \left(I - \frac{t}{g} A \right) \right) &= -2\sigma I + 2\sigma \left(-\frac{t}{g^3} A \mathbf{x}_T \mathbf{x}_T^T A^T + \frac{t}{g} A \right) \\ &= -2\sigma \left(I - \frac{t}{g} A \left(I - \frac{1}{g^2} \mathbf{x}_T \mathbf{x}_T^T A^T \right) \right). \end{aligned}$$

312 Hence, using Sylvester's Determinant Theorem, it follows that

$$(3.21) \quad \begin{aligned} -\frac{1}{(2\sigma)^3} \det \left(D \left(-2\sigma \mathbf{x}_T^T \left(I - \frac{t}{g} A \right) \right) \right) &= \det \left(I - \frac{t}{g} A \left(I - \frac{1}{g^2} \mathbf{x}_T \mathbf{x}_T^T A^T \right) \right) \\ &= \det \left(I - \frac{t}{g} \left(I - \frac{1}{g^2} \mathbf{x}_T \mathbf{x}_T^T A^T \right) A \right) \\ &= \det \left(I - \frac{t}{g} \left(I - \frac{1}{g^2} \mathbf{x}_T \mathbf{x}_T^T \right) A \right) \\ &= \det(I + C) \\ &= 1 + \text{tr}(C) + \frac{1}{2} ((\text{tr}(C))^2 - \text{tr}(C^2)) + \det(C), \end{aligned}$$

313 where we use SDT in the second step to reverse the matrix multiplication order, and the
314 fact that A is symmetric idempotent in the third step to get $A^T A = A^2 = A$. Here $C =$
315 $-\frac{t}{g} (I - B) A$, where $B = \frac{1}{g^2} \mathbf{x}_T \mathbf{x}_T^T$.

316 We now simplify (3.21). First, we have the identities

$$(3.22) \quad \text{tr}(A) = \text{tr} \left((r_1^T, r_2^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \right) = \text{tr} \left(\begin{pmatrix} r_1 \\ r_2 \end{pmatrix} (r_1^T, r_2^T) \right) = \text{tr}(I_{2 \times 2}) = 2,$$

317

$$(3.23) \quad \text{tr}(AB) = \frac{1}{g^2} \text{tr}(\mathbf{x}_T \mathbf{x}_T^T A) = \frac{1}{g^2} \text{tr}(\mathbf{x}_T^T A \mathbf{x}_T) = \frac{\mathbf{x}_T^T A \mathbf{x}_T}{g^2} = 1,$$

noting that $\mathbf{x}_T^T A \mathbf{x}_T = g^2$,

$$\text{tr}(BABA) = \frac{1}{g^4} \text{tr}(\mathbf{x}_T (\mathbf{x}_T^T A \mathbf{x}_T) \mathbf{x}_T^T A) = \frac{\mathbf{x}_T^T A \mathbf{x}_T}{g^4} \text{tr}(\mathbf{x}_T \mathbf{x}_T^T A) = \frac{(\mathbf{x}_T^T A \mathbf{x}_T)^2}{g^4} = 1,$$

and

$$\text{tr}(ABA) = \text{tr}(BA^2) = \text{tr}(BA) = 1.$$

Now

$$\operatorname{tr}(C) = -\frac{t}{g}\operatorname{tr}(A - BA) = -\frac{t}{g}(\operatorname{tr}(A) - \operatorname{tr}(BA)) = -\frac{t}{g},$$

318 and

$$\begin{aligned} \frac{1}{2}(\operatorname{tr}(C))^2 - \operatorname{tr}(C^2) &= \frac{1}{2} \left[\operatorname{tr}(C)^2 - \frac{t^2}{g^2} (\operatorname{tr}(BABA) - \operatorname{tr}(ABA) - \operatorname{tr}(BAA) + \operatorname{tr}(A^2)) \right] \\ &= \frac{t^2}{2g^2} [1 - (1 - 1 - 1 + 2)] = 0 \end{aligned}$$

and

$$\det(C) = -\frac{t^3}{g^3}\det(A)\det(I - B) = 0,$$

since $\det(A) = 0$. Indeed $0 \neq r_3^T \in \operatorname{Null}(A)$. Putting this together, we have

$$\det \left(D \left(-2\sigma \mathbf{x}_T^T \left(I - \frac{t}{g}A \right) \right) \right) = -(2\sigma)^3 \left(1 - \frac{t}{g} \right),$$

319 which is zero if and only if $t = g$. Hence, in the case when $t \neq g$, $D\Pi_L^{(1)}$ has full rank and
320 $\Pi_L^{(1)}$ is an immersion.

321 We now consider the case when $t = g$. In this case, $\mathbf{x} = R\mathbf{x}_C + \mathbf{x}_0$, where $\mathbf{x}_C =$
322 $(t \cos \theta, t \sin \theta, z)^T$, for some $\theta \in [0, 2\pi]$ and $z \in \mathbb{R}$. That is, \mathbf{x} lies on the cylinder of ra-
323 dius t , with axis of revolution $\{\mathbf{x}_0 + pR\mathbf{e}_3 : p \in \mathbb{R}\}$.

324 Under the assumption $t = g$, we show that the submatrix

$$(3.24) \quad M = \begin{array}{c} \text{d}_t \Phi_1 \\ \text{d}_\alpha \Phi_1 \\ s \end{array} \begin{array}{c} \nabla_{\mathbf{x}} \\ \nabla_{\mathbf{x}} \left(\frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T \right)^T \\ \nabla_{\mathbf{x}} (h - 2tg)^T \end{array}$$

of $D\Pi_L^{(1)}$ is invertible. Using the product rule and

$$\nabla_{\mathbf{x}} (g^i) = \frac{i}{2} g^{i-2} \nabla_{\mathbf{x}} (x'^2 + y'^2) = i g^{i-2} A \mathbf{x}_T,$$

325 for $i \in \mathbb{Z}$, we have

$$(3.25) \quad M = \begin{array}{c} \text{d}_t \Phi_1 \\ \text{d}_\alpha \Phi_1 \\ s \end{array} \begin{array}{c} -\frac{2\sigma}{g} (A \mathbf{x}_T)^T \\ 2\sigma \left[\frac{2t}{g} (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T - \frac{t}{g^3} \left(\mathbf{x}_T^T (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T \right) A \mathbf{x}_T \right]^T \\ 2 \left(\mathbf{x}_T - \frac{t}{g} A \mathbf{x}_T \right)^T \end{array}$$

326 Substituting $\mathbf{x} = R\mathbf{x}_C + \mathbf{x}_0$ and $t = g$, we have

$$(3.26) \quad M = \begin{array}{c} -\frac{2\sigma}{t} [(r_1^T, r_2^T, 0) \mathbf{x}_C]^T \\ 2\sigma \left[2 (r_{1\alpha}^T, r_{2\alpha}^T, 0) \mathbf{x}_C + \frac{1}{t^2} (tz \sin \beta \cos \theta) (r_1^T, r_2^T, 0) \mathbf{x}_C \right]^T \\ 2 (R\mathbf{x}_C - (r_1^T, r_2^T, 0) \mathbf{x}_C)^T \end{array}$$

327 where

$$\begin{aligned}
\mathbf{x}_T^T (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T &= \mathbf{x}_C^T R^T (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} R \mathbf{x}_C \\
&= \mathbf{x}_C^T R^T (r_{1\alpha}^T, r_{2\alpha}^T, 0) \mathbf{x}_C \\
(3.27) \quad &= \mathbf{x}_C^T \begin{pmatrix} 0 & 0 & 0 \\ \cos \beta & 0 & 0 \\ -\sin \beta & 0 & 0 \end{pmatrix} \mathbf{x}_C \\
&= -tz \sin \beta \cos \theta.
\end{aligned}$$

328 Here (3.27) shows the calculations for the new scalar term in brackets on the second row of
329 M in (3.26). We have

$$(3.28) \quad MR = \begin{pmatrix} -2\sigma(\cos \theta, \sin \theta, 0) \\ 2\sigma [2(-t \cos \beta \sin \theta, t \cos \beta \cos \theta, -t \sin \beta \cos \theta) + z \sin \beta \cos \theta(\cos \theta, \sin \theta, 0)] \\ 2(0, 0, z) \end{pmatrix}.$$

330 Thus

$$\begin{aligned}
\det(M) &= \det(MR) \\
&= 8\sigma^2 \begin{pmatrix} 0 \\ 0 \\ z \end{pmatrix} \cdot \left[- \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix} \times \left[2 \begin{pmatrix} -t \cos \beta \sin \theta \\ t \cos \beta \cos \theta \\ -t \sin \beta \cos \theta \end{pmatrix} + z \sin \beta \cos \theta \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix} \right] \right] \\
(3.29) \quad &= -16\sigma^2 z t \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \cdot \left[\begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix} \times \begin{pmatrix} -\cos \beta \sin \theta \\ \cos \beta \cos \theta \\ -\sin \beta \cos \theta \end{pmatrix} \right] \\
&= -16\sigma^2 z t \cos \beta.
\end{aligned}$$

331 Recall that $\beta \in (0, \pi/2)$ by definition of Y' , and $t > 0$. The case $z = 0$ corresponds to $s = 0$,
332 i.e., degenerate tori which have radius zero and collapse into a circle of radius t passing
333 through the center of the apple tube. We do not consider degenerate tori. Hence M , and
334 thus $D\Pi_L^{(1)}$, have full rank and $\Pi_L^{(1)}$ is an immersion.

335 3.2.2. $\Pi_L^{(1)}$ *injectivity proof*. Injectivity of $\Pi_L^{(1)}$ is a local property in the target space. To
336 determine if $\Pi_L^{(1)}$ is injective, we take an arbitrary point $(\mathbf{y}, \eta) \in T^*(\tilde{Y})$ and see if it has
337 more than one preimage. Specifically, we choose $\mathbf{y} \in \tilde{Y}$ and take local coordinates (3.11) so
338 that \mathbf{y} is in the image of Y' (i.e., the axis of the spindle torus parametrized by \mathbf{y} is neither
339 vertical nor horizontal). Then, we analyze $\Pi_L^{(1)}$ using these coordinates.

Let $\mathbf{x}_1, \mathbf{x}_2 \in B$ be such that

$$\Pi_L^{(1)}(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}_1; \sigma) = \Pi_L^{(1)}(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}_2; \sigma).$$

Then

$$2\sigma(t - g_1) = 2\sigma(t - g_2) \implies g_1 = g_2 = g,$$

where g_1, g_2 correspond to the inputs $\mathbf{x}_1, \mathbf{x}_2$. We now consider two cases, namely $t = g$ and
 $t \neq g$.

Case 1: $t \neq g$

We have

$$\left(I - \frac{t}{g}A\right)(\mathbf{x}_1 - \mathbf{x}_0) = \left(I - \frac{t}{g}A\right)(\mathbf{x}_2 - \mathbf{x}_0).$$

340 Thus, $\Pi_L^{(1)}$ is injective if $\left(I - \frac{t}{g}A\right)$ is invertible. Following similar arguments to those used
 341 in Theorem 3.2, we have

$$\begin{aligned} \det\left(I - \frac{t}{g}A\right) &= \det\left(I - \frac{t}{g}\begin{pmatrix} r_1^T & r_2^T \\ r_1 & r_2 \end{pmatrix}\begin{pmatrix} r_1 \\ r_2 \end{pmatrix}\right) \\ &= \det\left(I_{2 \times 2} - \frac{t}{g}\begin{pmatrix} r_1 \\ r_2 \end{pmatrix}\begin{pmatrix} r_1^T & r_2^T \end{pmatrix}\right) \\ (3.30) \quad &= \det\left(I_{2 \times 2} - \frac{t}{g}I_{2 \times 2}\right) \\ &= \left(1 - \frac{t}{g}\right)^2 \neq 0. \end{aligned}$$

342 Therefore, $\mathbf{x}_1 = \mathbf{x}_2$ and $\Pi_L^{(1)}$ is injective. Note we have used SDT in the second step of (3.30)
 343 to reverse the matrix multiplication order inside the determinant.

344

Case 2: $t = g$

346 In this case, $\mathbf{x}_j = R\mathbf{x}_C^{(j)} + \mathbf{x}_0$, for $j = 1, 2$, where $\mathbf{x}_C^{(j)} = (t \cos \theta_j, t \sin \theta_j, z_j)^T$ and $\theta_j \in [0, 2\pi]$.

347 That is, the \mathbf{x}_j lie on the cylinder, radius t , with axis of revolution $R\mathbf{e}_3 + \mathbf{x}_0$.

348 Using (3.12) and (3.13), we have

$$\begin{aligned} \nabla_{\mathbf{x}_0}\Phi_1(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}_j; \sigma) &= -2\sigma \left[R\mathbf{x}_C^{(j)} - AR\mathbf{x}_C^{(j)} \right] \\ (3.31) \quad &= -2\sigma \left[(r_1^T, r_2^T, r_3^T)\mathbf{x}_C^{(j)} - (r_1^T, r_2^T, 0)\mathbf{x}_C^{(j)} \right] \\ &= -2\sigma z_j r_3^T. \end{aligned}$$

349 Hence $z_1 = z_2 = z \neq 0$, since $r_3^T \neq \mathbf{0}$, and we do not consider the case $z = \sqrt{s} = 0$ (i.e.,
 350 a degenerate torus). Note, the \mathbf{x}_j are constrained also to lie on the apple parameterized by
 351 $(s, t, \mathbf{x}_0, \alpha, \beta)$, which, in the $t = g$ case, implies $z_j = \sqrt{s}$. Now,

$$\begin{aligned} d_\alpha\Phi_1(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}_j; \sigma) &= 2\sigma(\mathbf{x}_C^{(j)})^T R^T (r_{1\alpha}^T, r_{2\alpha}^T, 0) \mathbf{x}_C^{(j)} \\ (3.32) \quad &= 2\sigma(\mathbf{x}_C^{(j)})^T \begin{pmatrix} 0 & -\cos \beta & 0 \\ \cos \beta & 0 & 0 \\ -\sin \beta & 0 & 0 \end{pmatrix} \mathbf{x}_C^{(j)} \\ &= -2\sigma t z \cos \theta_j \sin \beta, \end{aligned}$$

352 and

$$\begin{aligned}
(3.33) \quad d_\beta \Phi_1(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}_j; \sigma) &= 2\sigma(\mathbf{x}_C^{(j)})^T R^T (r_{1\beta}^T, r_{2\beta}^T, 0) \mathbf{x}_C^{(j)} \\
&= 2\sigma(\mathbf{x}_C^{(j)})^T \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix} \mathbf{x}_C^{(j)} \\
&= -2\sigma t z \sin \theta_j.
\end{aligned}$$

353 Therefore, $(\cos \theta_1, \sin \theta_1) = (\cos \theta_2, \sin \theta_2)$ and $\Pi_L^{(j)}$ is injective. Recall that $\beta \in (0, \pi/2)$,
354 because $(s, t, \mathbf{x}_0, \alpha, \beta) \in Y'$, and so $\sin \beta > 0$.

355

356 This completes the proof of Theorem 3.3.

357 **3.3. Microlocal properties of \mathcal{L} ; the $j = 2$ case.** Here we discuss the microlocal prop-
358 erties of \mathcal{L} in a similar way to the $j = 1$ case. First, we prove that \mathcal{L} is an elliptic FIO order
359 -2 .

360 **Theorem 3.4.** *The lemon transform $\mathcal{L} = \mathcal{R}_2$ of (3.8) is an elliptic FIO order -2 from*
361 *domain $\mathcal{E}'(B)$ to $\mathcal{D}'(Y)$.*

362 *Proof.* As in Theorem 3.2, we choose local coordinates (3.11) on \tilde{Y} and use these local
363 coordinates in our calculations.

From (3.8), the phase function of \mathcal{R}_2 is

$$\Phi_2(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}; \sigma) = \sigma \Psi_2(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}).$$

364 We now show that Φ_2 is clean, non-degenerate and homogeneous in σ order 1, to show that
365 \mathcal{L} satisfies the definition of FIO (see definition 2.5). Φ_2 is trivially homogeneous order 1.
366 $d_s \Phi_2 = -\sigma \neq 0$, and hence $d\Phi_2, d_{\mathbf{y}} \Phi_2 \neq 0$. The lemon surfaces are smooth manifolds away
367 from their singular points, which we do not consider by the definition of Y , (3.6). Hence Φ_2
368 is clean.

369 Using similar calculations to those of (3.12), we have

$$\begin{aligned}
(3.34) \quad \nabla_{\mathbf{x}_0} \Phi_2 &= -2\sigma \left(I + \frac{t}{g} A \right) \mathbf{x}_T \\
&= -\sigma \nabla_{\mathbf{x}} \Psi_2 \\
&= -d_{\mathbf{x}} \Phi_2.
\end{aligned}$$

370 Also

$$\begin{aligned}
(3.35) \quad \det \left(I + \frac{t}{g} A \right) &= \det \left(I + \frac{t}{g} (r_1^T, r_2^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \right) \\
&= \det \left(I_{2 \times 2} + \frac{t}{g} \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} (r_1^T, r_2^T) \right), \text{ by (SDT)} \\
&= \det \left(I_{2 \times 2} + \frac{t}{g} I_{2 \times 2} \right) \\
&= \left(1 + \frac{t}{g} \right)^2 > 0.
\end{aligned}$$

371 Thus $d_{\mathbf{x}} \Phi_2$ is zero if and only if $\mathbf{x}_T = 0$, which we do not consider. Hence Φ_2 is nondegenerate.

372 The amplitude is

$$(3.36) \quad \begin{aligned} a_2(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}) &= \|\nabla_{\mathbf{x}} \Psi_2(s, t, \mathbf{x}_0, \alpha, \beta; \mathbf{x})\| \\ &= 2 \left(\mathbf{x}_T^T \left(I + \frac{t}{g} A \right)^T \left(I + \frac{t}{g} A \right) \mathbf{x}_T \right)^{\frac{1}{2}}. \end{aligned}$$

373 a_2 is smooth, and independent of σ , and hence a_2 is a symbol order zero. $a_2 > 0$ since $I + \frac{t}{g} A$
 374 is invertible, and hence $\left(I + \frac{t}{g} A \right)^T \left(I + \frac{t}{g} A \right)$ is positive definite. Therefore \mathcal{L} is an elliptic
 375 FIO order $O(\mathcal{L}) = 0 + \frac{1}{2} - \frac{7}{2} = -2$.

376 Recall that the spindle tori in \tilde{Y} do not have singular points in \bar{B} . Therefore, g is never
 377 zero and the symbol is defined for functions are supported in B . \square

378 We now have our second main theorem which shows that \mathcal{L} satisfies the semiglobal Bolker
 379 condition.

380 **Theorem 3.5.** *The left projection $\Pi_L^{(2)}$ of \mathcal{L} is an injective immersion, and hence \mathcal{L} satisfies*
 381 *the semiglobal Bolker condition.*

382 We now proceed in a similar fashion to the proof of Theorem 3.5, i.e., we split the proof
 383 into two subsections. We start with the immersion proof in the next section, and prove
 384 injectivity in the following section.

385 3.3.1. $\Pi_L^{(2)}$ immersion proof. We choose a point $(\mathbf{y}, \eta, \mathbf{x}, \xi)$ in the canonical relation of \mathcal{L} and
 386 choose coordinates as for the apple transform in section 3.2.1 so the spindle torus axis is
 387 neither vertical nor horizontal.

388 The left projection of \mathcal{R}_2 is

$$(3.37) \quad \begin{aligned} \Pi_L^{(2)}(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}; \sigma) &= \left(\overbrace{h + 2tg}^s, t, \mathbf{x}_0, \alpha, \beta, \overbrace{-\sigma}^{d_s \Phi_2}, \overbrace{-2\sigma \mathbf{x}_T^T \left(I + \frac{t}{g} A \right)}^{\nabla_{\mathbf{x}_0} \Phi_2} \right), \\ &\quad \underbrace{2\sigma(t+g)}_{d_t \Phi_2}, \underbrace{-\frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\alpha}^T, r_{2\alpha}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T}_{d_\alpha \Phi_2}, \underbrace{-\frac{2\sigma t}{g} \mathbf{x}_T^T (r_{1\beta}^T, r_{2\beta}^T) \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} \mathbf{x}_T}_{d_\beta \Phi_2}. \end{aligned}$$

389 The proof is analogous to the $j = 1$ case, and a little easier, so we will go over the main
 390 points.

We calculate $D\Pi_L^{(2)}$ and just consider the rows corresponding to $D_{\mathbf{x}}(\nabla_{\mathbf{x}_0} \Phi_2)$. We can
 show, in a similar way to the $j = 1$ case,

$$\det \left(D \left(-2\sigma \mathbf{x}_T^T \left(I + \frac{t}{g} A \right) \right) \right) = -(2\sigma)^3 \left(1 + \frac{t}{g} \right),$$

391 which is never zero. Therefore, these rows of $D\Pi_L^{(2)}$ have full rank 3. Hence $D\Pi_L^{(2)}$ has full
 392 rank and $\Pi_L^{(2)}$ is an immersion.

393 3.3.2. $\Pi_L^{(2)}$ *injectivity proof.* To prove injectivity we proceed similarly to Theorem 3.2, i.e., we
 394 take an arbitrary point $(\mathbf{y}, \eta) \in T^*(\tilde{Y})$ and determine whether it has more than one preimage
 395 under $\Pi_L^{(2)}$. We choose coordinates on \mathbb{R}^3 so that the axis of the lemon parameterized by \mathbf{y}
 396 is neither vertical nor horizontal.

Let $\mathbf{x}_1, \mathbf{x}_2 \in B$ be such that

$$\Pi_L^{(2)}(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}_1; \sigma) = \Pi_L^{(2)}(t, \mathbf{x}_0, \alpha, \beta; \mathbf{x}_2; \sigma) = (\mathbf{y}, \eta).$$

Then

$$2\sigma(t + g_1) = 2\sigma(t + g_2) \implies g_1 = g_2 = g,$$

397 where g_1, g_2 correspond to the inputs $\mathbf{x}_1, \mathbf{x}_2$.

Focusing on the $\nabla_{\mathbf{x}_0}\Phi_2$ terms in the image of $\Pi_L^{(2)}$ (see (3.37)), we have

$$\left(I + \frac{t}{g}A\right)(\mathbf{x}_1 - \mathbf{x}_0) = \left(I + \frac{t}{g}A\right)(\mathbf{x}_2 - \mathbf{x}_0).$$

398 Thus, $\Pi_L^{(2)}$ is injective if $\left(I + \frac{t}{g}A\right)$ is invertible. We have

$$(3.38) \quad \det\left(I + \frac{t}{g}A\right) = \left(1 + \frac{t}{g}\right)^2 > 0.$$

399 Hence, $\Pi_L^{(2)}$ is injective.

400

401 This concludes the proof of Theorem 3.5.

Corollary 3.6. *Let $\alpha \in [0, 2\pi]$, and $\beta \in [0, \pi/2]$ be fixed. Let $f \in \mathcal{E}'(B)$ and (s, t, \mathbf{x}_0) chosen so the singular points of the spindle torus parameterized by $(s, t, \mathbf{x}_0, \alpha, \beta)$ are disjoint from \bar{B} . Then the Radon transform*

$$\mathcal{L}_T f(s, t, \mathbf{x}_0) = \mathcal{L} f(s, t, \mathbf{x}_0, \alpha, \beta),$$

402 *which defines the integrals of f over a 5-D set of translated lemons, satisfies the semiglobal*
 403 *Bolker condition.*

The analogous restriction for the apple transform

$$\mathcal{A}_T f(s, t, \mathbf{x}_0) = \mathcal{A} f(s, t, \mathbf{x}_0, \alpha, \beta),$$

404 *however, does not satisfy the semiglobal Bolker condition.*

405 *Proof.* This follows immediately from Theorems 3.3 and 3.5. The left projection of \mathcal{A}_T
 406 has Jacobian which drops rank on the cylinder $t = g$, and thus there are artifacts in the
 407 reconstruction which occur along rings which are the intersections of apples and cylinders,
 408 radius t , with the same axis of revolution.

409 Regarding \mathcal{L} , we require only a 3-D translation of the lemons, and the radial variables (s
 410 and t), in order for the semiglobal Bolker condition to be satisfied. The rotations induced
 411 by α, β are not needed in the proof of Theorem 3.5. \square

412 3.3.3. *Discussion.* In [24], lemon transforms are analyzed, but only rotations and changes in
 413 radius of the lemons are considered (i.e., $\mathbf{x}_0 = 0$ is fixed, and α, β , and s, t vary). The authors
 414 prove that the left projection drops rank, and show that there are artifacts in (unfiltered)
 415 backprojection, and Landweber image reconstructions. With knowledge of seven-dimensional
 416 lemon integral data, however, we would not expect to see artifacts due to rank deficiencies
 417 in the reconstruction. In fact, five-dimensional lemon integral data is sufficient to show the
 418 Bolker condition is satisfied, as is shown by Corollary 3.6.

419 With regards to \mathcal{A} , the full seven-dimensional data is needed in the proof of Theorem 3.3
 420 to show that the Bolker condition is satisfied. As noted in Corollary 3.6, there is issue with
 421 the translated apples on their intersections with cylinders, radius t , which share the same
 422 axis of revolution. Such issues can be addressed by including the 2-D rotation induced by α
 423 and β . Thus, with knowledge of seven-dimensional apple integral data, we would not expect
 424 to see artifacts due to microlocal properties in the reconstruction.

425

426 In the following section, we consider 3-D subsets of apple and lemon surfaces which have
 427 practical motivations in CST. In the case of the apple transform, we discover artifacts which
 428 occur at apple-cylinder intersections, and are thus consistent with the results of Theorem
 429 3.3.

430

4. PRACTICAL GEOMETRY IN CST

431 In this section we consider the machine geometry of [23, 27], which has practical applica-
 432 tions in airport baggage screening. We present a microlocal analysis of the apple transform,
 433 first introduced in [23], and its lemon transform analog. Specifically, we consider the machine
 geometry of figure 4.

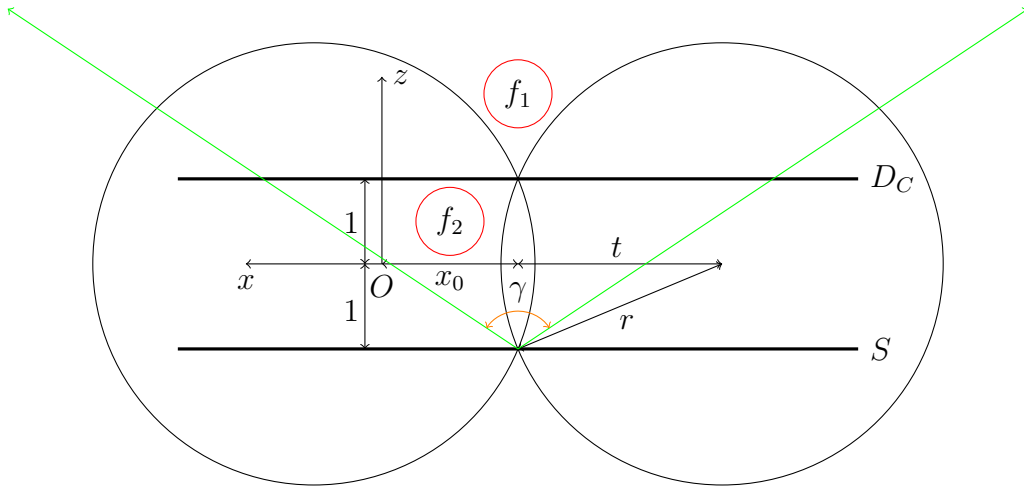


FIGURE 4. Two-dimensional cross-section of parallel CST geometry. S and D_C denote the source and detector rows, which are parallel and one unit away from the origin. The y axis is perpendicular to the page. Here we show a 2-D cross section of a spindle torus with center \mathbf{x}_0 , and $t > 0$ is the distance from \mathbf{x}_0 to the center of the torus tube, as before. The torus has radius $r = \sqrt{t^2 + 1} > 0$.

434

435 The diagram illustrates an X-ray scanner comprised of a line segment of sources (S), which
436 emit X-rays in the direction of a parallel line segment of detectors (D_C). The photons are
437 then Compton scattered and measured by the detectors on D_C . Meanwhile, the target, f ,
438 is translated out of the page (i.e., in the y axis direction) on a conveyor belt. We consider
439 two possibilities for the location of f here, namely within the half space $\{z > 1\}$ (i.e., above
440 D_C), and within the band $\{-1 < z < 1\}$ (i.e. between S and D_C). Examples of these
441 two possible locations for f are illustrated by f_1 and f_2 in figure 4, respectively, where f_1 is
442 integrated over apples and f_2 over lemons. The source is cone beam with opening angle γ .
443 We set $\gamma = \pi$ so that photons are everywhere on $\{z > -1\}$. See [23, 27], for more details on
444 the applications to airport baggage screening and CST, more generally. In total, the data is
445 three-dimensional, and is comprised of a 2-D translation and a 1-D radial variable. In this
446 section, $R = I$ is fixed (i.e., there is no rotation of the apples or lemons), $z_0 = 0$ (i.e., the
447 translation is in the (x, y) plane), and s and t satisfy the relation $s = t^2 + 1$ (see figure 4).
448 With this in mind we define the restricted apple and lemon transforms

$$(4.1) \quad \begin{aligned} \mathcal{A}_0 f(p, x_0, y_0) &= \mathcal{A}f \left(\frac{p}{4} + 1, \sqrt{\frac{p}{4}}, (x_0, y_0, 0)^T, 0, 0 \right), \\ &\text{and} \\ \mathcal{L}_0 f(p, x_0, y_0) &= \mathcal{L}f \left(\frac{p}{4} + 1, \sqrt{\frac{p}{4}}, (x_0, y_0, 0)^T, 0, 0 \right), \\ &\text{where } p = 4t^2. \end{aligned}$$

449 The variable p is introduced in this section to simplify the calculations.

450 **Proposition 4.1.** *The restricted apple and lemon transforms can be written*

$$(4.2) \quad \mathcal{A}_0 f = \mathcal{T}_1 f \text{ for } f \in L_c^2(\{z > 1\}), \quad \mathcal{L}_0 f = \mathcal{T}_2 f, \text{ for } f \in L_c^2(\{-1 < z < 1\})$$

451 where

$$(4.3) \quad \begin{aligned} \mathcal{T}_j f(p, x_0, y_0) &= \int_{X_j} \|\nabla_{\mathbf{x}} \Psi(p, x_0, y_0; \mathbf{x})\| \delta(\Psi(p, x_0, y_0; \mathbf{x})) f(\mathbf{x}) d\mathbf{x} \\ &= \int_{-\infty}^{\infty} \int_{X_j} \|\nabla_{\mathbf{x}} \Psi(p, x_0, y_0; \mathbf{x})\| e^{\sigma \Psi(p, x_0, y_0; \mathbf{x})} f(\mathbf{x}) d\mathbf{x} d\sigma, \end{aligned}$$

for $j = 1, 2$, where we now define

$$(4.4) \quad \begin{aligned} X_1 &= \{\mathbf{x} \in \mathbb{R}^3 : z > 1\}, \quad X_2 = \{\mathbf{x} \in \mathbb{R}^3 : (x - x_0)^2 + (y - y_0)^2 + z^2 < 1\}, \\ \Psi(p, x_0, y_0; \mathbf{x}) &= p - \frac{h^2}{g} \quad \text{where} \\ g &= g(x_0, y_0; \mathbf{x}) = (x - x_0)^2 + (y - y_0)^2, \quad h = h(x_0, y_0; \mathbf{x}) = \|\mathbf{x}_T\|^2 - 1, \\ \mathbf{x}_T &= (x - x_0, y - y_0, z) \end{aligned}$$

452 Note that the functions Ψ, g, h, \mathbf{x}_T are adapted from section 3 for our geometry.

453 *Proof.* A torus centered at the origin with axis of rotation z is described implicitly by the
454 equation

$$(4.5) \quad (\|\mathbf{x}\|^2 + t^2 - s)^2 = 4t^2(x^2 + y^2).$$

455 Hence, the defining equation for the tori of interest, which are translated by (x_0, y_0) in the
 456 (x, y) plane (as depicted in figure 4), and satisfy $t^2 - s = -1$, becomes

$$(4.6) \quad \Psi(p, x_0, y_0; \mathbf{x}) = p - \frac{((x - x_0)^2 + (y - y_0)^2 + z^2 - 1)^2}{(x - x_0)^2 + (y - y_0)^2} = p - \frac{h^2(x_0, y_0; \mathbf{x})}{g(x_0, y_0; \mathbf{x})}$$

457 Thus, when the integration is restricted to X_1 , $\mathcal{T}_1 f$ defines the integrals of f over a 3-D set of
 458 translated apples whose singular points lie on $\{z = 1\}$ and $\{z = -1\}$. $\mathcal{T}_2 f$ defines integrals
 459 of f over lemons in the same way when the integration is restricted to X_2 .

460 Note that the functions in $L_c^2(\{z > 1\})$ and $L_c^2(\{-1 < z < 1\})$ the domains of \mathcal{A}_0 and
 461 \mathcal{L}_0 , respectively, in (4.2) are zero near the singular points of the spindle tori (which satisfy
 462 $z = \pm 1$). Hence the surface measure on spindle tori for \mathcal{T}_j is defined on the support of f . \square

463 We now show that the \mathcal{T}_j are elliptic FIO order -1 .

464 **Theorem 4.2.** *The Radon transforms \mathcal{T}_j , for $j = 1, 2$, are elliptic FIO order -1 from
 465 domain $\mathcal{E}'(\{z > 1\})$ for $j = 1$ and from domain $\mathcal{E}'(\{-1 < z < 1\})$ for $j = 2$.*

466 *Proof.* For the proofs in this section, it will be convenient to define the function

$$(4.7) \quad u(x_0, y_0; \mathbf{x}) = \frac{h(x_0, y_0; \mathbf{x})}{g(x_0, y_0; \mathbf{x})}.$$

467 The phase function of \mathcal{T}_j , for $j = 1, 2$, is

$$(4.8) \quad \Phi(p, x_0, y_0; \mathbf{x}; \sigma) = \sigma \Psi(p, x_0, y_0; \mathbf{x}) = \sigma(p - hu),$$

468 by (4.3), and the amplitude is

$$(4.9) \quad a(p, x_0, y_0; \mathbf{x}) = \|\nabla_{\mathbf{x}} \Psi(p, x_0, y_0; \mathbf{x})\| = 2u\sqrt{g(u - 2)^2 + 4z^2}.$$

469 The phase (4.8) and the amplitude (4.9) are undefined when $g = 0$, that is on the rotation
 470 axis of each spindle torus—when $(x, y) = (x_0, y_0)$. To get around this, we use a smooth cutoff
 471 near the spindle torus axis as in [26, Lemma 3.3] to smoothly set the symbol to zero near
 472 the spindle torus axis. Note that the points at which the cutoff is not smooth, the singular
 473 points of the spindle torus (on $z = \pm 1$), are not in either domain in (4.2). This cutoff
 474 makes the amplitude defined and smooth everywhere. Note that the phase is smooth on
 475 a neighborhood of the canonical relations of our transforms and the cutoff on the symbol
 476 can be used to make it smooth everywhere. Similarly, for \mathcal{L}_0 , the integral in (4.3) for each
 477 (x_0, y_0) is over the open disk, X_2 so as to integrate only over the lemon, not the part of the
 478 apple in $\{-1 < z < 1\}$. One constructs a smooth function of $(p, x_0, y_0; \mathbf{x})$ that is equal to 1
 479 in a neighborhood in $(0, \infty) \times \mathbb{R}^2 \times \{-1 < z < 1\}$ of the lemon parameterized by (p, x_0, y_0)
 480 and equal to 0 in a neighborhood of the corresponding apple.

481 The phase (4.8) is trivially clean and homogeneous in σ order 1. In addition, $d_p \Phi = \sigma \neq 0$,
 482 and

$$\|d_{\mathbf{x}} \Phi\| = 2\sigma u \sqrt{g(u - 2)^2 + 4z^2}.$$

$u = 0$ does not occur since $p > 0$. When $j = 1$, the domain of integration is such that $z > 1$,
 and hence \mathcal{T}_2 has nondegenerate phase. When $j = 2$, $-1 < z < 1$, and

$$u - 2 = \frac{1}{g} [z^2 - (x - x_0)^2 - (y - y_0)^2 - 1] < 0.$$

483 Hence \mathcal{T}_1 has nondegenerate phase.

484 By the same arguments as for the phase, the amplitude (4.9) is positive on the apple and
 485 on the lemon, the manifolds of integration of \mathcal{A}_0 and \mathcal{L}_0 . By using the cutoff, a is smooth.
 486 Furthermore, a does not depend on σ , so a is an elliptic symbol order zero, and thus \mathcal{T}_j , for
 487 $j = 1, 2$, is an elliptic FIO order $O(\mathcal{T}_j) = 0 + \frac{1}{2} - \frac{3}{2} = -1$. \square

488 We now have our third main theorem which provides conditions such that the \mathcal{T}_j satisfy
 489 the semiglobal Bolker condition.

490 **Theorem 4.3.** *Global coordinates on the canonical relation of \mathcal{T}_j are given by $(x_0, y_0; \mathbf{x}; \sigma)$,*
 491 *as in (3.11) but with $(\alpha, \beta) = (0, 0)$ and $p = h^2(x_0, y_0; \mathbf{x})/g(x_0, y_0; \mathbf{x})$. Let*

$$(4.10) \quad D_1 = \mathbb{R}^2 \times \{\mathbf{x} \in \mathbb{R}^3 : z > 1\} \times \mathbb{R} \setminus \{0\}$$

492 and

$$(4.11) \quad D_2 = \{(x_0, y_0; \mathbf{x}) \in \mathbb{R}^5 : z \in (0, 1)\} \times \mathbb{R} \setminus \{0\}.$$

493 *These sets define global coordinates on the appropriate canonical relation.*

494 *The left projection $\Pi_L^{(1)} : D_1 \rightarrow \Pi_L^{(1)}(D_1)$ of \mathcal{T}_1 is an injective immersion under the con-*
 495 *straint that*

$$(4.12) \quad u - 2 > 0, \quad \text{equivalently } z^2 - 1 > (x - x_0)^2 - (y - y_0)^2.$$

496 *The left projection $\Pi_L^{(2)} : D_2 \rightarrow \Pi_L^{(2)}(D_2)$ of \mathcal{T}_2 is an injective immersion.*

497 *Remark 4.4.* The requirement in Theorem 4.3 that functions are supported in the half-space
 498 $z > 0$ (or equivalently $z < 0$) is natural because the apples and lemons are symmetric about
 499 $z = 0$. Therefore, our transforms integrate odd functions in z to zero and singularities for
 500 $z < 0$ can cancel singularities for $z > 0$.

501 We point out that (4.12) puts restrictions on the support of functions and the sets of
 502 (x_0, y_0) for which \mathcal{A}_0 satisfies the Bolker condition. Define the set

$$(4.13) \quad H(x_0, y_0) = \{(x, y, z) : z > 1, z^2 - 1 > (x - x_0)^2 - (y - y_0)^2\}.$$

503 Let $f \in L_c^2(\{z > 1\})$. If $\text{supp}(f)$ is so large that it is not contained in $H(x_0, y_0)$ for any
 504 (x_0, y_0) , then one cannot apply Theorem 4.3 to f .

505 Now, assume $(x_0, y_0) \in \mathbb{R}^2$ and K is a compact subset of $z > 1$ such that $K \subset H(x_0, y_0)$.
 506 Then by compactness of K , there is an open neighborhood, U , of (x_0, y_0) such that $K \subset$
 507 $H(x_1, y_1)$ for all $(x_1, y_1) \in U$. Theorem 4.3 can be applied to the local problem for \mathcal{A}_0 for
 508 functions supported in K and centers (x_0, y_0) in U .

509 *Proof.* The points $(x_0, y_0; \mathbf{x}; \sigma)$ are coordinates on the canonical relation of \mathcal{T}_j for the same
 510 reason as (3.11) give local coordinates on \tilde{Y} . However, here they are global coordinates
 511 because $(\alpha, \beta) = (0, 0)$ is fixed.

512 Let $u_1 = u - 1$ and $u_2 = u - 2$. Then, the left projection of \mathcal{T}_j , for $j = 1, 2$, is

$$(4.14) \quad \Pi_L^{(j)}(\sigma; x_0, y_0; \mathbf{x}) = \left(\underbrace{d_p \Phi}_{\sigma}, x_0, y_0, \underbrace{p}_{uh}, \underbrace{d_{x_0} \Phi}_{-2\sigma u u_2(x - x_0)}, \underbrace{d_{y_0} \Phi}_{-2\sigma u u_2(y - y_0)} \right),$$

513 where we have rearranged the variables to correspond to the order used in calculating the
 514 Jacobian

$$(4.15) \quad D\Pi_L^{(j)} = \begin{pmatrix} I_{3 \times 3} & 0_{3 \times 3} \\ \cdot & M_2 \end{pmatrix},$$

where, using

$$d_x(uu_2) = 2u_x u_1 = -\frac{4}{g}(x - x_0)u_1^2,$$

and

$$d_y(uu_2) = 2u_y u_1 = -\frac{4}{g}(y - y_0)u_1^2,$$

515 we have

$$(4.16) \quad M_2 = \begin{matrix} p \\ d_{x_0}\Phi \\ d_{y_0}\Phi \end{matrix} \begin{pmatrix} \begin{matrix} dx & & \\ -2u_2(x - x_0)u & \frac{8\sigma}{g}(x - x_0)(y - y_0)u_1^2 & \frac{-8\sigma}{g}(x - x_0)zu_1 \\ \frac{8\sigma}{g}(x - x_0)(y - y_0)u_1^2 & -2\sigma(-4(y - y_0)^2u_1^2 + hu_2) & -\frac{8\sigma}{g}(y - y_0)zu_1 \end{matrix} \\ \begin{matrix} dy & & \\ -2u_2(y - y_0)u & & 4zu \\ & & dz \end{matrix} \end{pmatrix}.$$

516 The determinant of M_2 is hence

$$(4.17) \quad \det(M_2) = \frac{16z\sigma^2}{g^2}u \times \det(M_3),$$

517 where

$$(4.18) \quad M_3 = \begin{pmatrix} -(x - x_0)u_2 & -(y - y_0)u_2 & 1 \\ 4(x - x_0)^2u_1^2 - hu_2 & 4(x - x_0)(y - y_0)u_1^2 & -2(x - x_0)u_1 \\ 4(x - x_0)(y - y_0)u_1^2 & 4(y - y_0)^2u_1^2 - hu_2 & -2(y - y_0)u_1 \end{pmatrix}.$$

518 A straightforward calculation shows that $\det(M_3) = -h^2uu_2$, and hence

$$(4.19) \quad \det D\Pi_L^{(j)} = -16z\sigma^2u^4(u - 2).$$

519 Therefore, $\det D\Pi_L^{(j)} = 0$ if and only if $z = 0$, $u = 0$ or $u = 2$. On D_1 , $u, z > 0$ and hence
 520 the left projection of \mathcal{T}_1 drops rank if and only if $u = 2$. By condition (4.12), $u > 2$, and
 521 hence $\det D\Pi_L^{(1)}$ is nonzero and $\Pi_L^{(1)}$ is an immersion. On D_2 , $u_2, u < 0$ and hence the left
 522 projection of \mathcal{T}_2 drops rank if and only if $z = 0$, which we do not consider by assumption
 523 that $z > 0$. Hence $\det D\Pi_L^{(2)}$ is an immersion.

Now onto injectivity. First, we consider the case $j = 1$. Let $\mathbf{x}_1 = (x_1, y_1, z_1)$ and $\mathbf{x}_2 = (x_2, y_2, z_2)$ be such that

$$\Pi_L^{(1)}(x_0, y_0; \mathbf{x}_1; \sigma) = \Pi_L^{(1)}(x_0, y_0; \mathbf{x}_2; \sigma),$$

524 and let $v = u(x_0, y_0; \mathbf{x}_1)$, and $w = u(x_0, y_0; \mathbf{x}_2)$. Let $v_1 = v - 1$, $v_2 = v - 2$, $w_1 = w - 1$, and
 525 $w_2 = w - 2$.

$$(4.20) \quad (vh_1, (x_1 - x_0)vv_2, (y_1 - y_0)vv_2) = (wh_2, (x_2 - x_0)ww_2, (y_2 - y_0)ww_2),$$

526 where $h_j = h(x_0, y_0; \mathbf{x}_j)$, and $g_j = g(x_0, y_0; \mathbf{x}_j)$, for $j = 1, 2$. It follows that

$$\begin{aligned}
 & [(x_1 - x_0)^2 + (y_1 - y_0)^2] v^2 v_2^2 = [(x_2 - x_0)^2 + (y_2 - y_0)^2] w^2 w_2^2 \\
 \implies & g_1 v^2 v_2^2 = g_2 w^2 w_2^2 \\
 \implies & v h_1 v_2^2 = w h_2 w_2^2 \\
 \implies & v_2^2 = w_2^2, \quad (\text{note } v h_1 = w h_2 = p > 0).
 \end{aligned}
 \tag{4.21}$$

527 Under our assumption that $v_2, w_2 > 0$, it follows that $v_2 = w_2$, so $v = w$ and $h_1 = h_2$. Using
 528 (4.20) again, we see that $v = w \implies (x_1, y_1) = (x_2, y_2)$ (note $v, w \neq 0, 2$), and so $z_1^2 = z_2^2$
 529 since $h_1 = h_2$. On D_1 , $z > 1$, so $z_1 = z_2$ and $\Pi_L^{(1)}$ is thus an injective immersion.

530 On D_2 , $v_2, w_2 < 0$, and hence by (4.21) and the previous arguments $\Pi_L^{(2)}(x_0, y_0; \mathbf{x}_1; \sigma) =$
 531 $\Pi_L^{(2)}(x_0, y_0; \mathbf{x}_2; \sigma)$ implies that $(x_1, y_1) = (x_2, y_2)$ and $z_1^2 = z_2^2$. Also, $z \in (0, 1)$, which implies
 532 $z_1 = z_2$, and thus $\Pi_L^{(2)}$ is an injective immersion. \square

533 **4.1. Discussion of the artifacts.** In this section, we discuss the restrictions imposed on
 534 the function support and left projection domain in Theorem 4.3 needed to show that the
 535 semiglobal Bolker condition is satisfied, and address the artifacts that occur when such
 constrains are lifted.

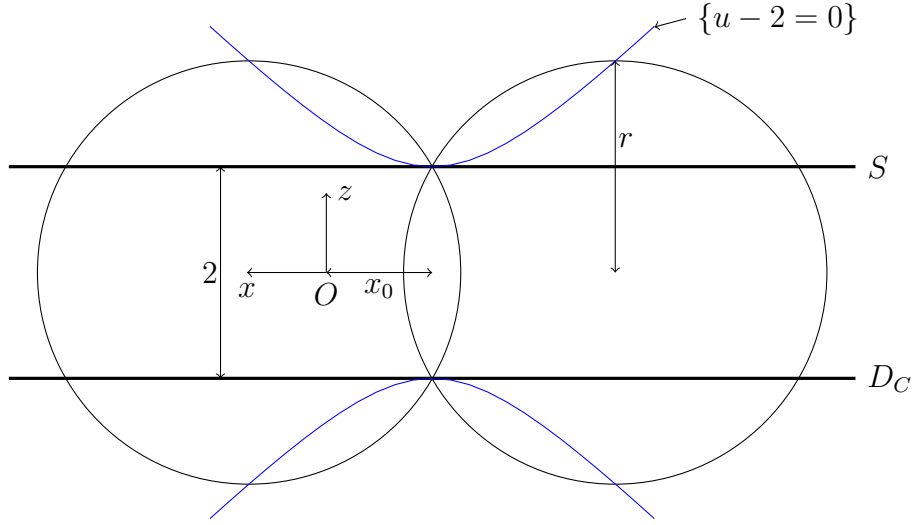


FIGURE 5. Intersection of a torus and $\{u - 2 = 0\}$ shown as a 2-D cross-section. A 2-D cross-section of $\{u - 2 = 0\}$ is drawn in blue and intersects the translated torus at $z = \pm r$.

536

537 In the proof of Theorem 4.3, the constraint $z \in (0, 1)$ was needed to show that $\Pi_L^{(2)}$ is an
 538 injective immersion. Without loss of generality, we could replace the $z \in (0, 1)$ constraint
 539 with $z \in (-1, 0)$ and the proof would follow in the same way, by symmetry of the apples and
 540 lemons about $z = 0$. If the function support is not restricted in this way, and z takes values on
 541 the full range $-1 < z < 1$, then $\Pi_L^{(2)}$ is noninjective. For example, $\Pi_L^{(2)}((x, y, z); \sigma, x_0, y_0) =$
 542 $\Pi_L^{(2)}((x, y, -z); \sigma, x_0, y_0)$, for any $-1 < z < 1$, and thus there are artifacts which consist
 543 of reflections in the (x, y) plane. When $z = 0$, specifically, $D\Pi_L^{(2)}$ drops rank and $\Pi_L^{(2)}$

544 is nonimmersive. $\Pi_L^{(1)}$ also suffers the same noninjectivity concerns if z takes both signs.
 545 However, it does not practically make sense for the function to be supported on both sides
 546 of the (x, y) plane when integrating over apples. Indeed, the cone-beam direction in figure
 547 4 is such that there are no photons on $\{z < -1\}$, and thus we assume f is supported on
 548 $\{z > 1\}$ when integrating over apples (as is done also in [27]).

549 Regarding $\Pi_L^{(1)}$, $u > 2$ was assumed a-priori in Theorem 4.3, and discussed also in Remark
 550 4.4, in order to show that $\Pi_L^{(1)}$ satisfies the Bolker condition. Without such restrictions,
 551 in particular when $u = 2$, $D\Pi_L^{(1)}$ drops rank and there artifacts which occur along rings at
 552 the top and bottom of the apple surface. Specifically, when $u = 2$, \mathbf{x} lies on the two-sided
 553 hyperboloid, described implicitly by

$$(4.22) \quad z^2 - (x - x_0)^2 - (y - y_0)^2 - 1 = 0.$$

554 The intersection of the apple and the surface defined by (4.22) occurs when $z = \pm\sqrt{t^2 + 1} =$
 555 r , i.e., along the rings at the top and bottom of the apple. See figure 5, where we have shown
 556 a 2-D cross section of the intersecting apple and hyperboloid surfaces. In Corollary 3.6,
 557 we showed that \mathcal{A}_T did not satisfy the Bolker condition. Specifically, the left projection
 558 of \mathcal{A}_T drops rank for \mathbf{x} on $\{t = g\}$, namely the cylinders radius t , with axis of revolution
 559 $\{\mathbf{x}_0 + \nu R\mathbf{e}_3 : \nu \in \mathbb{R}\}$, i.e., the axis of revolution of the apple surface. The apple and $\{t = g\}$
 560 intersect on rings at the top and bottom of the apple which are the same intersection points
 561 as those shown in figure 5. Thus, our results are consistent with the findings of Corollary
 562 3.6. Specifically, when the degrees of freedom in our data includes translation, e.g., the full
 563 3-D translation of $\mathcal{A}_T f$, or the 2-D translation of $\mathcal{A}_0 f$, there is a consistency in the artifact
 564 locations.

565 **4.2. How to remove artifacts with machine design.** In this section, we discuss possible
 566 modifications to the machine design of figure 4, so that the conditions of Theorem 4.3 are
 567 met, and thus we do not have to contend with the types of artifacts discussed in the previous
 568 section.

569 When using forward scattered photons for imaging, whose intensity is modeled by the
 570 lemon transform, we need only restrict the support of f to $\{0 < z < 1\}$ or $\{-1 < z < 0\}$.
 571 See figure 6 for an example f with such support, in particular the location of f_2 . In this
 572 case, the conditions of Theorem 4.3 are satisfied and the lemon transform satisfies the Bolker
 573 condition. Practically speaking, such support restrictions can be achieved by re-positioning
 574 the scanning target (e.g., the airport luggage) to be strictly above or below the (x, y) plane.
 575 For example, we could construct the conveyor belt to lie on $\{z = 0\}$ (highlighted by a red
 576 dashed line in figure 6) and place the scanning target (with height less than 1) on top of the
 577 conveyor, to ensure the conditions of Theorem 4.3 are met.

578 Regarding \mathcal{A}_0 , i.e., when backscattered photons are used for imaging, the object is com-
 579 pactly supported on $\{z > 1\}$. To ensure the conditions of Theorem 4.3 are met, we propose
 580 to further restrict the support of f to $\{z > 1 + \epsilon\}$, for some $\epsilon > 0$. See f_1 in figure 6 for
 581 an example f with such support. In practice, this would mean placing the conveyor belt on
 582 $\{z = 1 + \epsilon\}$ (shown as a red dashed line in figure 6), with f on top of the conveyor. With
 583 such restrictions on the support of f , we can choose the cone-beam angle γ (as shown in
 584 figure 6) so that no scatter occurs on the surface $\{u - 2 = 0\}$, and $u - 2 > 0$. Note that we
 585 have removed the bottom half of $\{u - 2 = 0\}$ in figure 6, since f_1 is supported on $\{z > 1 + \epsilon\}$.

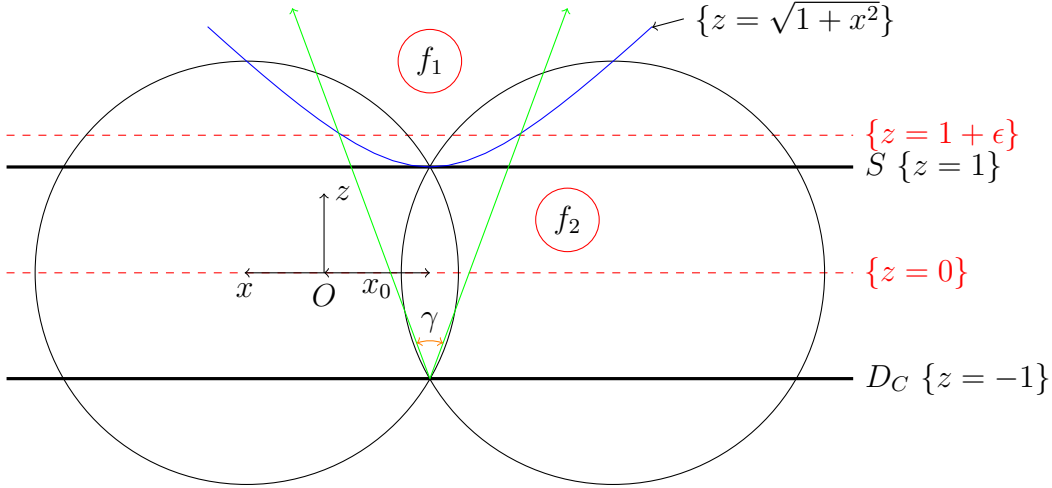


FIGURE 6. Restrictions to source cone beam angle (backscatter) and object support (forward scatter) so that Bolker is satisfied.

586 To restrict the scatter exclusively to $u - 2 > 0$, we can write γ explicitly as

$$(4.23) \quad \gamma = \min \left(2 \tan^{-1} \frac{\sqrt{(1+\epsilon)^2 - 1}}{2 + \epsilon}, 90^\circ \right).$$

587 Note, γ must be less than or equal to 90° since if $\gamma > 90^\circ$, the line through $(x_0, -1)$ and
588 $(x_0 \pm \sqrt{(1+\epsilon)^2 - 1}, 1 + \epsilon)$ (the right/left hand green line of figure 6) has gradient less than
589 1, and hence intersects the blue curve of figure 6 for large enough x (i.e., scatter could occur
590 on or below $\{u - 2 = 0\}$). Note that the blue curve in figure 6 (i.e., $\{z = \sqrt{1+x^2}\}$) has
591 max gradient 1, so the green lines of figure 6 must have gradient greater than or equal to
592 1 to ensure that $u - 2 > 0$ and the conditions of Theorem 4.3 are satisfied. In practice,
593 such restrictions on γ would mean there is less signal, due to the smaller cone-beam and
594 less photons, and hence the data would become more noisy. So, while we can address the
595 microlocal artifacts by restricting γ , this would in turn increase the noise level. Thus, there
596 is a trade of to consider here, i.e., do we want higher Signal to Noise Ratio (SNR) and more
597 artifacts, or less SNR with less artifacts? We leave such practical concerns for future work.

598 5. CONCLUSIONS AND FURTHER WORK

599 In this paper, we presented a novel microlocal analysis of seven-dimensional apple and
600 lemon Radon transforms, which have applications in CST. The goal of this work was to
601 consider a best case scenario in CST, in terms of data dimensionality. The literature [24,
602 25, 18, 23, 16, 17, 2] considers exclusively Radon transforms which define the integrals of
603 a function over three-dimensional sets of apple or lemon surfaces. In these works, artifacts
604 are present in the reconstruction due to data limitations, and regularization strategies are
605 used to combat the artifacts. Here, we considered a case when a full seven-dimensional set
606 of apple and lemon integrals are known. Our main theorems, namely Theorems 3.3 and 3.5),
607 prove that the apple and lemon transforms are elliptic FIO, order 2, which satisfy the Bolker
608 condition.

609 In addition, we investigated the microlocal properties of apple and lemon transforms which
610 induce translation of the target function, and discussed an example machine geometry from
611 airport baggage screening, first introduced in [27]. We analyzed two lemon transforms,
612 namely \mathcal{L}_T and \mathcal{L}_0 (see Corollary 3.6 and (4.1) respectively), which were shown to satisfy the
613 Bolker condition when the function support was restricted to the upper half of the unit ball.
614 The corresponding apple transforms \mathcal{A}_T and \mathcal{A}_0 were shown to violate the Bolker condition.
615 Specifically, there were artifacts induced on the intersections of apples and cylinders with
616 the same axis of revolution. This indicates higher instability in \mathcal{A}_T and \mathcal{A}_0 inversion, when
617 compared to \mathcal{L}_T and \mathcal{L}_0 . Thus, it may be beneficial to use forward scattered photons,
618 which correspond to lemon integrals, if one were to manufacture a CST machine with linear
619 scanning motion (e.g., a scanner with translated sources and detectors, as considered here).
620 The theory is not global however, and does not account for all CST geometries which include
621 linear motion. In further work, we aim to generalize our theory and determine whether such
622 artifacts as discovered here are present for any CST modality with translated sources and
623 detectors.

624 REFERENCES

- 625 [1] L. Borg, J. Frikel, J. S. Jørgensen, and E. T. Quinto. Analyzing reconstruction artifacts from arbitrary
626 incomplete X-ray CT data. *SIAM Journal on Imaging Sciences*, 11(4):2786–2814, 2018.
- 627 [2] J. Cebeiro, C. Tarpau, M. A. Morvidone, D. Rubio, and M. K. Nguyen. On a three-dimensional compton
628 scattering tomography system with fixed source. *Inverse Problems*, 37(5):054001, 2021.
- 629 [3] J. J. Duistermaat. *Fourier integral operators*, volume 130 of *Progress in Mathematics*. Birkhäuser, Inc.,
630 Boston, MA, 1996.
- 631 [4] J. J. Duistermaat and L. Hörmander. *Fourier integral operators*, volume 2. Springer, 1996.
- 632 [5] V. Guillemin and S. Sternberg. *Geometric Asymptotics*. American Mathematical Society, Providence,
633 RI, 1977.
- 634 [6] L. Hörmander. *The analysis of linear partial differential operators. I*. Classics in Mathematics. Springer-
635 Verlag, Berlin, 2003. Distribution theory and Fourier analysis, Reprint of the second (1990) edition
636 [Springer, Berlin].
- 637 [7] L. Hörmander. *The analysis of linear partial differential operators. III*. Classics in Mathematics.
638 Springer, Berlin, 2007. Pseudo-differential operators, Reprint of the 1994 edition.
- 639 [8] L. Hörmander. *The analysis of linear partial differential operators. IV*. Classics in Mathematics.
640 Springer-Verlag, Berlin, 2009. Fourier integral operators, Reprint of the 1994 edition.
- 641 [9] C.-Y. Jung and S. Moon. Inversion formulas for cone transforms arising in application of Compton
642 cameras. *Inverse Problems*, 31(1):015006, 2015.
- 643 [10] A. I. Katsevich. Local tomography for the limited-angle problem. *J. Math. Anal. Appl.*, 213(1):160–182,
644 1997.
- 645 [11] V. P. Krishnan and E. T. Quinto. Microlocal analysis in tomography. *Handbook of mathematical methods*
646 *in imaging*, pages 1–50, 2014.
- 647 [12] P. Kuchment and F. Terzioglu. Three-dimensional image reconstruction from Compton camera data.
648 *SIAM Journal on Imaging Sciences*, 9(4):1708–1725, 2016.
- 649 [13] S. Moon and M. Haltmeier. Analytic inversion of a conical radon transform arising in application of
650 Compton cameras on the cylinder. *SIAM Journal on imaging sciences*, 10(2):535–557, 2017.
- 651 [14] M. K. Nguyen, T. T. Truong, and P. Grangeat. Radon transforms on a class of cones with fixed axis
652 direction. *Journal of Physics A: Mathematical and General*, 38(37):8003, 2005.
- 653 [15] E. T. Quinto. The dependence of the generalized Radon transform on defining measures. *Trans. Amer.*
654 *Math. Soc.*, 257:331–346, 1980.
- 655 [16] G. Rigaud. 3d compton scattering imaging with multiple scattering: Analysis by fio and contour recon-
656 struction. *Inverse Problems*, 2021.

- 657 [17] G. Rigaud and B. Hahn. Reconstruction algorithm for 3d compton scattering imaging with incomplete
658 data. *Inverse Problems in Science and Engineering*, 29(7):967–989, 2021.
- 659 [18] G. Rigaud and B. N. Hahn. 3D Compton scattering imaging and contour reconstruction for a class of
660 Radon transforms. *Inverse Problems*, 34(7):075004, 2018.
- 661 [19] W. Rudin. *Functional analysis*. McGraw-Hill Book Co., New York, 1973. McGraw-Hill Series in Higher
662 Mathematics.
- 663 [20] Sylvester’s determinant theorem. [http://www.scientificlib.com/en/Mathematics/LX/
664 SylvestersDeterminantTheorem.html](http://www.scientificlib.com/en/Mathematics/LX/SylvestersDeterminantTheorem.html), 2014. Accessed 11/30/2021.
- 665 [21] J. J. Sylvester. On the relation between the minor determinants of linearly equivalent quadratic func-
666 tions. *Philosophical Magazine*, 1:295–305, 1851.
- 667 [22] T. T. Truong, M. K. Nguyen, and H. Zaidi. The mathematical foundations of 3D Compton scatter
668 emission imaging. *International journal of biomedical imaging*, 2007, 2007.
- 669 [23] J. Webber and E. Miller. Compton scattering tomography in translational geometries. Technical report,
670 Tufts University, 2019.
- 671 [24] J. W. Webber and S. Holman. Microlocal analysis of a spindle transform. *Inverse Problems & Imaging*,
672 13(2):231–261, 2019.
- 673 [25] J. W. Webber and W. R. Lionheart. Three dimensional Compton scattering tomography. *Inverse Prob-
674 lems*, 34(8):084001, 2018.
- 675 [26] J. W. Webber and E. T. Quinto. Microlocal analysis of generalized radon transforms from scattering
676 tomography. *SIAM Journal on Imaging Sciences*, 14(3):976–1003, 2021.
- 677 [27] J. W. Webber, E. T. Quinto, and E. L. Miller. A joint reconstruction and lambda tomography regular-
678 ization technique for energy-resolved x-ray imaging. *Inverse Problems*, 36(7):074002, 2020.